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The activation gate controls steady-state inactivation and recovery from inactivation in *Shaker*

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Despite major advances in the structure determination of ion channels, the sequence of molecular rearrangements at negative membrane potentials in voltage-gated potassium channels of the *Shaker* family remains unknown. Four major composite gating states are documented during the gating process: closed (C), open (O), open-inactivated (OI), and closed-inactivated (CI). Although many steps in the gating cycle have been clarified experimentally, the development of steady-state inactivation at negative membrane potentials and mandatory gating transitions for recovery from inactivation have not been elucidated. In this study, we exploit the biophysical properties of *Shaker*-IR mutants T449A/V474C and T449A/V476C to evaluate the status of the activation and inactivation gates during steady-state inactivation and upon locking the channel open with intracellular Cd²⁺. We conclude that at negative membrane potentials, the gating scheme of *Shaker* channels can be refined in two aspects. First, the most likely pathway for the development of steady-state inactivation is C→O→OI⇌CI. Second, the OI→CI transition is a prerequisite for recovery from inactivation. These findings are in accordance with the widely accepted view that tight coupling is present between the activation and C-type inactivation gates in *Shaker* and underscore the role of steady-state inactivation and recovery from inactivation as determinants of excitability.

Introduction

C-type inactivation is a process that renders voltage-gated potassium (K_v) channels nonconductive upon prolonged depolarization in the absence of the fast, N-type inactivation. N-type inactivation is mediated by the block of the permeation pathway by the N-terminal inactivation particle, whereas C-type inactivation depends on the amino acids in and around the selectivity filter close to the C-termini of the channel subunits. C-type inactivation of potassium channels plays a major role in the governance of diverse biological processes such as modulating the firing of neurons (Bean, 2007), regulating the length and frequency of the action potential in cardiomyocytes (Smith et al., 1996; Jan and Jan, 2012; Bocksteins, 2016), and regulating the activation of nonexcitable immune (Matteson and Deutsch, 1984; Panyi et al., 2004; Villalonga et al., 2010; Pérez-García et al., 2018) and endocrine cells (MacDonald and Wheeler, 2003). C-type inactivation is widely prevalent across many K⁺ channel families (Hoshi et al., 1991; Kurata and Fedida, 2006; Cuello et al., 2010b; Hoshi and Armstrong, 2013). Structurally, C-type inactivation is linked to conformational changes at the selectivity filter (Kurata and Fedida, 2006; Cuello et al., 2010b; Hoshi and Armstrong, 2013; Pless et al., 2013; Pau et al., 2017; Matthies et al., 2018). Furthermore, novel studies emphasize the

importance of “inactivation water” molecules in the cavity behind the selectivity filter in controlling the molecular rearrangements related to inactivation (Cordero-Morales et al., 2011; Cuello et al., 2017) and recovery from inactivation (Ostmeyer et al., 2013).

The *Shaker*-IR channel lacking the N-type inactivation balls is a dedicated model for examining the C-type inactivation of K_v channels (Hoshi et al., 1990, 1991; Kurata and Fedida, 2006). In *Shaker*-IR, the flow of ions is controlled by the coupled opening and closing of the activation gate (A-gate), formed by the bundle-crossing of four S6 helices (Yellen, 1998; Bezanilla, 2000; del Camino and Yellen, 2001; Webster et al., 2004), and the C-type inactivation gate (I-gate) formed by the selectivity filter. Four major gating states can be defined: closed (C), open (O), open-inactivated (OI), and closed-inactivated (CI) states, as indicated in Fig. 1 (Yellen, 1998; Panyi and Deutsch, 2006; Ostmeyer et al., 2013; Cuello et al., 2017). The existence of all the gating states is confirmed by functional, structural, and computational studies (Long et al., 2005; Cuello et al., 2010a, 2010b; Pau et al., 2017; Labro et al., 2018; Matthies et al., 2018).

The opening and closing of the A-gate is controlled by the membrane potential. Membrane depolarization opens, while

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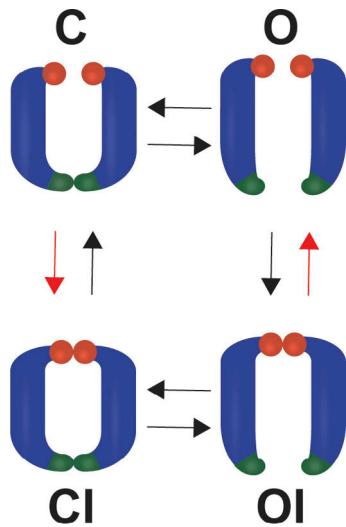


Figure 1. **The four-state gating model of the *Shaker*-IR channel.** The gating of *Shaker*-IR channels can be described by the four-state gating model, in which four major gating states can be distinguished depending on the status of the activation (dark green) and inactivation (red spheres) gates. Two opposing pore domains are illustrated in the cartoons. Among the closed (C), open (O), open-inactivated (OI), and closed-inactivated (CI) states, only the open state is conducting. The theoretically existing transitions between these gating states are indicated by arrows. In this study we investigate the existence of direct C→CI and OI→O transitions (red arrows) at negative membrane potentials.

repolarization closes, the A-gate (Yellen, 1998; Bezannila, 2000). Upon A-gate opening, extended depolarization leads to C-type inactivation (O→OI), while sufficiently negative membrane potential triggers closing of the A-gate (OI→CI) long before recovery from inactivation takes place (CI→C; Kurata and Fedida, 2006; Panyi and Deutsch, 2006). The A-gate and I-gate are bidirectionally coupled, mediated by allosteric communication trajectories between the gates (Ogielska and Aldrich, 1999; Panyi and Deutsch, 2006; Sadosky and Yifrach, 2007; Cuello et al., 2010a, 2010b; Labro et al., 2018).

Transitions at depolarized membrane potentials that lead to channel opening (O state) only partially determine the amount of current available for regulating physiological processes. It is equally important to understand what fraction of the channels is available to open. This is mainly determined by transitions at negative membrane potentials, i.e., the relative probability of channels in CI and C at steady state. The phenomenon of steady-state inactivation reflects this distribution. What determines this $CI \rightleftharpoons C$ equilibrium? What promotes recovery of closed I-gates? To address these issues, we consider putative gating transitions and ask the following questions. First, can the closed channel directly inactivate at negative membrane potentials (C→CI)? Second, must the inactivated channel (OI state) pass through the CI state to recover from inactivation, or can it bypass this state to recover? Can the inactivation gate (the I-gate), in its non-conducting configuration, rearrange/open while the A-gate is continuously open, i.e., is an OI→O transition permitted at negative voltages for recovery to occur? The answers to these mechanistic questions will expand our understanding of

K_V gating transitions and aid the design of more efficacious drugs with state-dependent binding, e.g., to sites in the cavity of the channel, which depend on the probability that the A-gate and/or the I-gate are open.

Studying gating transitions at negative membrane potentials is hampered by two factors. One is that at these membrane potentials, the open probability of the channels is very small, resulting in miniature currents. In addition, the driving force for K^+ is small, which also results in small currents and a large error in the determination of the state of the A-gate from current measurements. In contrast, gating transitions at the A-gate, i.e., movement of the A-gate, can be precisely tracked using the state-dependent cysteine accessibility method (Liu et al., 1997).

Both C→CI and OI→O transitions require the I-gate to rearrange without A-gate movements, which contradicts the primary coupling between the activation and C-type inactivation gates. Our aim was to investigate the existence of the direct C→CI and OI→O transitions at negative membrane potentials in *Shaker*-IR channels transiently expressed in human embryonic kidney (HEK293) cells. To this end, we employed the patch-clamp technique in inside-out configuration and T449A/V474C *Shaker*-IR channels for studying closed-state inactivation (direct C→CI transition) or T449A/V476C *Shaker*-IR channels to examine recovery from inactivation in the presence of a locked-open A-gate (direct OI→O transition). These electrophysiologically silent transitions were visualized using state-dependent cysteine accessibility, Cd^{2+} modification, and Cd^{2+} crosslinking.

Materials and methods

DNA clones and site-directed mutagenesis

The *Shaker*-IR construct in a GW1-CMV mammalian expression plasmid containing a deletion of amino acids 6–46 to remove N-type inactivation, and C301S and C308S point mutations (Holmgren et al., 1996) to exclude Cd^{2+} modification of endogenous *Shaker* cysteines and accompanying functional effects, were a kind gift from R. Horn (Thomas Jefferson University, Philadelphia, PA; Ding and Horn, 2002). Amino acid substitutions at positions 449 (T449A) and 474 (V474C) or 476 (V476C), according to *Shaker* B numbering, were introduced using a QuikChange site-directed mutagenesis kit (Stratagene).

Cell culture and transfection

The human embryonic kidney cell line, HEK293, was obtained from ATCC and grown according to the instructions of the manufacturer. A calcium phosphate transfection kit (Invitrogen) was used to cotransfect the EGFP plasmid with the mutant *Shaker* constructs using a ratio of 1:10 μ g. Transfected cells were replated into Corning 35-mm polystyrene cell culture dishes pretreated with poly-L-ornithine (Sigma-Aldrich) for higher cell adhesion to excise inside-out patches. Then channels were transiently expressed in the cells for 12–36 h after transfection. The EGFP-positive transfectants were identified on a Nikon TE2000U fluorescence microscope using bandpass filters of 455–495 nm and 515–555 nm for excitation and emission, respectively. In general, >60% of the EGFP-positive cells expressed the cotransfected ion channels.

Electrophysiology

All the experiments were performed using inside-out configuration. The direction of K^+ currents and the voltage protocols are presented according to general conventions. Data were acquired using an Axopatch 200B amplifier (Molecular Devices) and digitized with Axon Digidata 1550 (Molecular Devices). The filter frequency was set to less than or equal to half of the sampling frequency. Pipettes were pulled from GC 150 F-15 borosilicate capillaries (Harvard Apparatus), and resistances were 8–9 M Ω . Only patches with a steady-state current <5% of the peak current were used in these studies. All experiments were performed at room temperature (20–24°C).

Solutions and perfusion system

The standard intracellular (bath) solution contained (in mM) 105 KF, 35 KCl, 10 EGTA, and 10 HEPES, titrated to pH 7.36–7.38 with KOH, with a final K^+ concentration of 160–165 mM and osmolarity of 285–295 mOsm/L. To characterize the perfusion kinetics, the standard intracellular solution was identical, except that it contained 50 mM K^+ and 100 mM NaF instead of the corresponding K^+ salts and was titrated with NaOH. For experiments using Cd^{2+} as the modifying reagent, the intracellular solution was identical to the standard K^+ -based solution except that EGTA was omitted and (in mM) 125 KF, 35 KCl, and 10 HEPES were used. Standard extracellular (pipette) solution was (in mM) 150 NaCl, 2 KCl, 1.5 $CaCl_2$, 1 $MgCl_2$, and 10 HEPES. The osmolarity of the solution was 290 mOsm/L, and pH was titrated to 7.36–7.38 with NaOH.

A Warner Instruments SF-77A Perfusion Fast-Step system with 3-barrel square glass (700 μ m internal side length) was used for rapid solution exchange. Inside-out patches were perfused with standard intracellular solutions at a rate of 0.5 ml/min. The principle and the description of the test protocol to determine the rate of the solution exchange has been described previously (Panyi and Deutsch, 2006, 2007). The solution exchange was complete in 30 ms (see Results).

Data analysis

Data were analyzed using Clampfit (v10; Molecular Devices), SigmaPlot (v10.; Systat Software), and Excel (Microsoft). Before analysis, inside-out current traces were digitally filtered with a three-point boxcar filter. All the ionic currents applied for determining conductance–voltage (G – V) and steady-state inactivation curves were leak corrected before the analysis.

To construct normalized G (G_{norm})– V curves of mutant *Shaker*-IR channels, inside-out patches were held at –120 mV and depolarized for 100 ms to test potentials ranging from –100 to +70 mV in steps of 10 mV every 60 s. Conductances at given test potentials were obtained from the leak-corrected peak currents (I_{peak}) at test potential (V) and the K^+ reversal potential (E_{rev}) using $G = I_{\text{peak}}/(V - E_{\text{rev}})$. While isochronal tail current calculation is the standard method for determining G – V curves for noninactivating currents, the *Shaker*-IR/T449A channels in this study are inactivating, which complicates the interpretation of tail current analysis. The G values were normalized for the maximum conductance (G_{norm}) and plotted as a function of the test potential. G_{norm} – V curves were quantified by fitting a

Boltzmann function ($y = 1/[1 + e^{-(V-V_{1/2})/k}]$) and determining half-activation voltage ($V_{1/2}$) and slope factor (k) from the fits.

To determine the inactivation time constants (τ_{inact}) of the currents, 2,000-ms-long depolarizing pulses to +50 mV were applied. The decaying part of the current traces was fitted by a single exponential function ($I = I_0 \times [e^{-t/\tau}] + C$) starting from 90% of the peak. To define activation time constants (τ_{act}) of the currents, 5-ms-long depolarizing pulses were applied from a holding potential of –120 mV to +50 mV, then evoked current traces were fitted using Hodgkin-Huxley n^4 -model ($I = I_0[1 - e^{-t/\tau}]^4 + C$).

For measuring the kinetics of recovery from inactivation, pairs of depolarizing pulses were delivered from the holding potential of –120 to +50 mV for 200 ms. The interpulse interval (ipi) at –120 mV varied between 0.5 and 60 s. The fractional recovery was calculated as $(I_2 - I_{\text{SS}})/(I_1 - I_{\text{SS}})$, where I_2 and I_1 are the peak currents during the second and first pulse, respectively, and I_{SS} is the steady-state current at the end of the first depolarization. Data points were fitted by a single exponential function ($I = I_0 \times (1 - e^{-t/\tau}) + C$) to obtain the time constant of recovery from inactivation (τ_{rec}).

To describe steady-state inactivation curves, the fraction of noninactivated channels at each voltage was calculated as I/I_{-120} , where I is the peak current evoked by the depolarization from a given prepulse potential, and I_{-120} is the peak current evoked by identical depolarization from the holding potential of –120 mV. $V_{1/2}$ and k were determined by fitting a Boltzmann function to the data points.

Data are represented as mean \pm SEM. The numbers of cells (n) involved in the given analysis are shown in the text. P values were calculated based on ANOVA followed by Tukey's honestly significant difference test. Differences were considered significant at $P < 0.05$.

Online supplemental material

Fig. S1 shows raw current traces corresponding to the analysis in Fig. 4, B and C (holding potential of –120 mV and –90 mV between pulses P_1 and P_2 in A and B, respectively, no Cd^{2+} applied). Fig. S2 shows raw current traces corresponding to the analysis in Fig. 5, B and C (holding potential of –90 mV between pulses P_1 and P_2 in the presence of 200 μ M Cd^{2+}). Fig. S3 shows the biophysical characterization of T449A/V476C *Shaker*-IR channels in inside-out patches.

Results

Experimental strategy to investigate the pathway leading to inactivation at negative membrane potentials in the T449A/V476C *Shaker*-IR channel

In *Shaker*-IR channels, residue 474 is located in the S6 helix, facing the water-filled cavity of the pore when A-gate is open (Liu et al., 1997). A cysteine introduced into this position can only be modified by Cd^{2+} or MTS reagents in the case of an open A-gate (Liu et al., 1997; Panyi and Deutsch, 2006, 2007). Cd^{2+} binding blocks ion conduction. T449A amino acid substitution was employed to speed up C-type inactivation to keep the length of our protocols reasonable (López-Barneo et al., 1993; Hoshi and

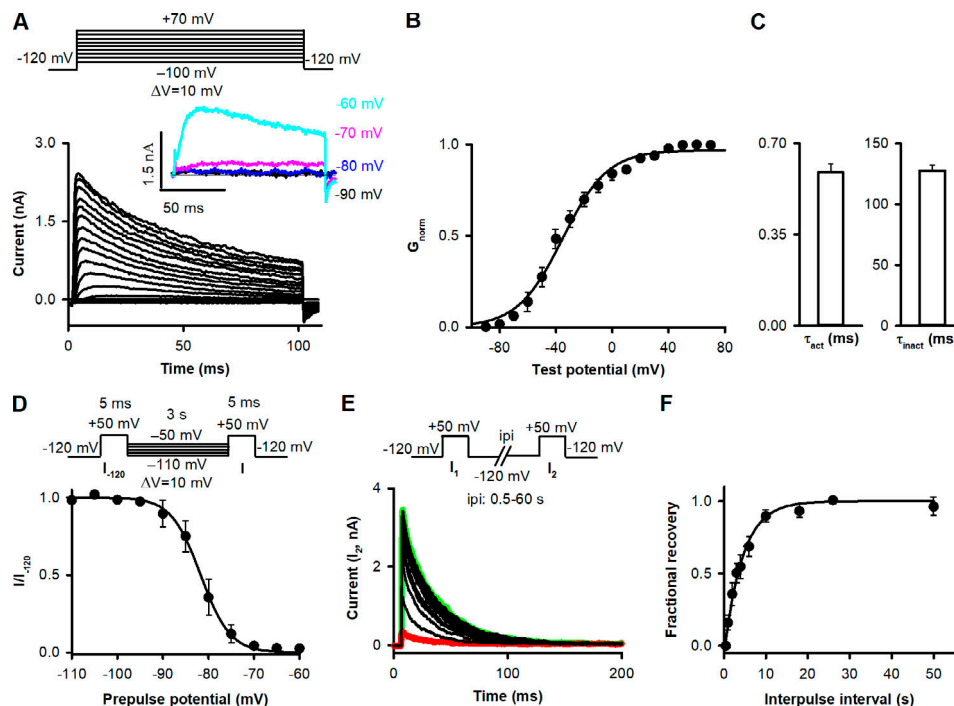


Figure 2. Electrophysiological parameters of T449A/V474C Shaker-IR ion channels expressed in HEK293 cells. (A) K^+ currents in inside-out patch configuration for the T449A/V474C Shaker-IR channel. The patch was held at -120 mV and depolarized for 100 ms to test potentials ranging from -100 to $+70$ mV in steps of 10 mV every 60 s (see protocol above). Inset shows currents evoked by the depolarizing pulses to the indicated membrane potentials. (B) G_{norm} -V data (see Materials and methods) were obtained for $n = 5$ cells, averaged, and plotted as a function of test potentials. The superimposed solid line shows the best-fit Boltzmann function to the averaged data points. (C) To determine activation time constant (τ_{act}), we applied 5-ms-long depolarizing pulses from a holding potential of -120 to $+50$ mV, and the evoked current traces were fitted using the Hodgkin-Huxley n^4 -model ($n = 6$). Inactivation time constant (τ_{inact}) of the current was determined by fitting a single exponential function to the decaying part of the currents evoked by 2,000-ms-long pulses to $+50$ mV ($n = 5$). (D) To describe voltage dependence of steady-state inactivation, patches were stepped from a holding of -120 mV to prepulse potentials between -110 and -50 mV in steps of 5 mV for 3 s, and then a short (5-ms) test pulse to $+50$ mV was applied to elicit potassium currents (I ; see protocol above). After each step, a second test pulse was applied from a holding potential of -120 mV (I_{-120}). The fraction of noninactivated channels (I/I_{-120}) at each voltage was averaged for $n = 4$ cells and plotted as a function of prepulse potential. The superimposed solid line shows the best-fit Boltzmann function to the averaged data points. (E) To describe recovery from inactivation, pairs of depolarizing pulses were delivered from the holding potential of -120 mV to $+50$ mV for 200 ms (see protocol above). Currents recorded during the second pulses (I_2) are shown at various iipi at -120 mV holding potential. Traces recorded with iipi of 0.5 and 60 s are highlighted in red and green, respectively. (F) The fractional recovery was calculated as $(I_2 - I_{\text{SS1}})/(I_1 - I_{\text{SS1}})$, where I_1 and I_2 are the peak currents during the first and second pulse, and I_{SS1} is the steady-state current at the end of the first depolarization. Fractional recovery values were averaged for $n = 5$ cells for each iipi and plotted as a function of iipi. The best-fit exponential function to the averaged data points is indicated in the figure. Error bars indicate SEM throughout the figure.

Armstrong, 2013). To design a pulse sequence for optimizing the study of state transitions and kinetics, we electrophysiologically characterized the T449A/V474C Shaker-IR channel using inside-out patch-clamp configuration (see Data analysis). First the current-voltage (Fig. 2 A) and voltage dependence of the conductance were determined (Fig. 2 B) by fitting a Boltzmann function to the normalized conductance (G_{norm})-test potential relationships to give $V_{1/2} = -51.6 \pm 2.5$ mV and a slope factor of $k = 16.9 \pm 2.1$ mV ($n = 5$). The activation and inactivation time constants of the currents at $+50$ mV are in Fig. 2 C. The current activates quickly, and the inactivation kinetics is fast enough to conduct the experiments on a reasonable time scale. The voltage dependence of steady-state inactivation was characterized by $V_{1/2}$ ($V_{1/2} = -81.7 \pm 0.7$ mV, $n = 4$) and a slope factor ($k = -5.8 \pm 0.6$ mV, $n = 4$) obtained from Boltzmann fits for I/I_{-120} dependence on prepulse potential of individual cells (see Materials and methods). The I/I_{-120} averaged at each prepulse potential is shown in Fig. 2 D. Using pulse pairs to $+50$ mV and varying the

iipi, the kinetics of recovery from inactivation were determined at -120 mV, which served as the holding potential for the experiments between pulse protocols (Fig. 2 E), to give a recovery time constant of 4.4 ± 0.6 s ($n = 5$; Fig. 2 F).

Based on these biophysical properties of T449A/V474C channels, we designed an experimental strategy using state-dependent cysteine accessibility to Cd^{2+} applied from the intracellular side to investigate closed-state inactivation (Fig. 3 A). A pulse sequence consisting of three 5-ms-long depolarizations to $+50$ mV was applied (Fig. 3 A). These depolarizing steps fully activated potassium currents (Fig. 2 B). During the first pulse (P_1), control solution (150 mM K^+ , no Cd^{2+}) was applied to obtain the maximum current (I_1). This was followed by a 30-ms-long hyperpolarization to -120 mV to ensure quick and full closure of the A-gate of all channels before Cd^{2+} application (Fig. 3 A, symbol \times). The membrane potential was then stepped to various negative values for 20 s. As a control, the membrane potential was stepped to -120 mV, where neither steady-state inactivation

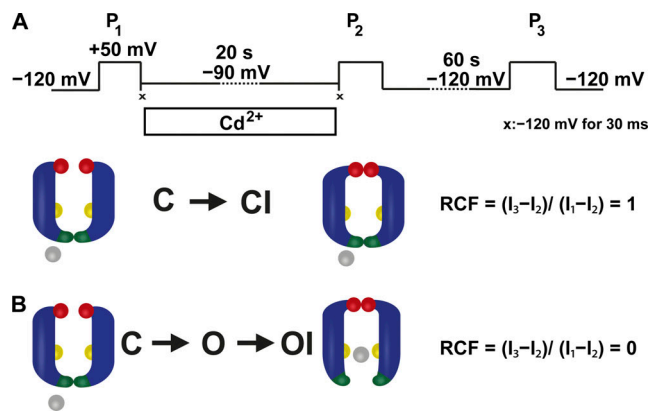


Figure 3. Experimental strategy to investigate closed-state inactivation in T449A/V474C Shaker-IR ion channels. (A) The pulse protocol consisted of three 5-ms-long depolarizations to +50 mV (P_1 , P_2 , and P_3). The corresponding peak currents evoked by these pulses are represented as I_1 , I_2 , and I_3 in B. Between P_1 and P_2 , patches were held in Cd^{2+} containing intracellular solution at -90 mV holding potential for 20 s. To avoid the Cd^{2+} exposure of the open-state channels during P_1 and P_2 , 30-ms-long steps to -120 mV were applied directly after P_1 and before P_2 , indicated by symbol \times . The holding potential of -120 mV was applied between P_2 and P_3 for 60 s. (B) The accessibility of cysteines at position 474 (yellow spheres) located between the activation (intracellular part of S6 helices, dark green) and C-type inactivation (red spheres) gates during inactivation was determined by applying Cd^{2+} (gray spheres) from the intracellular side. Upon binding to 474C, Cd^{2+} inhibits K^+ currents, thus tracking precisely the incidental opening of the A-gate during the development of C-type inactivation. Model predictions for RCF are in the text; I_1 , I_2 , and I_3 are defined in A.

nor A-gate opening was observed (Fig. 4). The status of the A-gate was then assessed at membrane potentials at which the channel shows a small degree of steady-state inactivation (Fig. 2 D), but no macroscopic current can be detected (Fig. 2, A and B, inset in A), e.g., at -90 mV (Fig. 3 A). Although Fig. 2 D may be an underestimate of true steady-state inactivation at -90 mV, it nonetheless defines the appropriate range of negative potentials for examination of state-dependent Cd^{2+} accessibility. The extent of steady-state inactivation at -90 mV was assessed by the current (I_2) evoked by the second depolarizing pulse (P_2). P_2 was also preceded by a 30-ms-long step to -120 mV to avoid Cd^{2+} exposure of the open-state channels at the beginning of P_2 (electromechanical delay in the perfusion system; Panyi and Deutsch, 2006; Fig. 3 A, symbol \times). In the absence of Cd^{2+} , the ratio of the peak currents (I_2/I_1) during P_2 and P_1 was used to quantify steady-state inactivation. Thereafter the patch was held at -120 mV for 60 s, which was sufficient to ensure full recovery of the inactivated channels in control conditions. Consequently, the current amplitude (I_3) during the third 5-ms-long depolarizing pulse (P_3) in control cases was similar to that of the first pulse. When examining the modification of 474C by Cd^{2+} , the pulse sequence was essentially the same as for control cases except the 20-s-long pulse was applied in the presence of an intracellular solution containing either 20 or 200 μM Cd^{2+} . Cd^{2+} exposure of patches was initiated after P_1 and terminated with the completion of the -90-mV holding step, before P_2 . The inactivated current fraction (IF) was calculated as $\text{IF} = 1 - (I_2/I_1)$, and the recovered current fraction as $\text{RCF} = (I_3 - I_2)/(I_1 - I_2)$

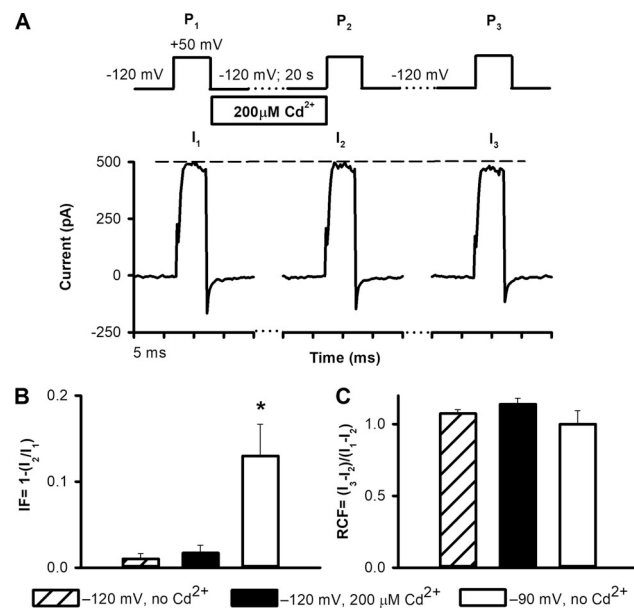


Figure 4. The three-pulse protocol: control measurements. Three different experimental conditions were used for control measurements on HEK293 cells expressing T449A/V474C Shaker-IR channels in inside-out patch. (A) Representative traces demonstrate currents evoked by the three depolarizing pulses to +50 mV (P_1 , P_2 , and P_3) using a holding of -120 mV. The intracellular solution perfusate contained 200 μM Cd^{2+} during the 20-s holding, as indicated. The dashed reference line is drawn to I_1 . (B) Inactivated fraction of the channels ($\text{IF} = [I_1 - I_2]/I_1$) at -120 mV holding between P_1 and P_2 (hatched bar), at -120 mV in the presence of 200 μM Cd^{2+} (solid bar), and at -90 mV in the absence of Cd^{2+} (empty bar). (C) $\text{RCF} = (I_3 - I_2)/(I_1 - I_2)$ at -120 mV holding between P_1 and P_2 (hatched bar), at -120 mV holding between P_1 and P_2 in the presence of 200 μM Cd^{2+} (solid bar), and at -90 mV holding between P_1 and P_2 in the absence of Cd^{2+} (empty bar). Data are represented as mean \pm SEM ($n = 5$) throughout the figure. Differences were considered significant (*) compared with control (-120 mV, no Cd^{2+}) when $P < 0.05$.

(Fig. 3 B). I_1 , I_2 , and I_3 represent the leak-corrected current amplitudes evoked by the three sequential pulses. Theoretically, if inactivation originates solely from the closed state, i.e., there is no opening of the A-gate before inactivation, then RCF should equal 1, because the closed A-gate prevents Cd^{2+} modification of 474C and current loss (Fig. 3 B, upper model). In contrast, if A-gate opening precedes inactivation, Cd^{2+} will modify 474C, and current loss will be seen. If all channels contributing to steady-state inactivation have an open A-gate at -90 mV and Cd^{2+} modifies all of these channels, then RCF is 0 (Fig. 3 B, lower model). $0 < \text{RCF} < 1$ indicates that the channels during steady-state inactivation were in a Cd^{2+} -accessible state (see Discussion).

Activation gate opening in the T449A/V474C Shaker-IR channel at negative holding potentials

To evaluate the ability of the three-pulse protocol to report specifically on the opening of the A-gate, we used three different control experiments. First, application of a -120-mV holding potential for 20 s between P_1 and P_2 in the absence of Cd^{2+} did not decrease I_2 compared with I_1 (Fig. S1 A). No inactivation occurred ($\text{IF} \sim 0$) and consequently, RCF is close to 1 (Fig. 4, B and

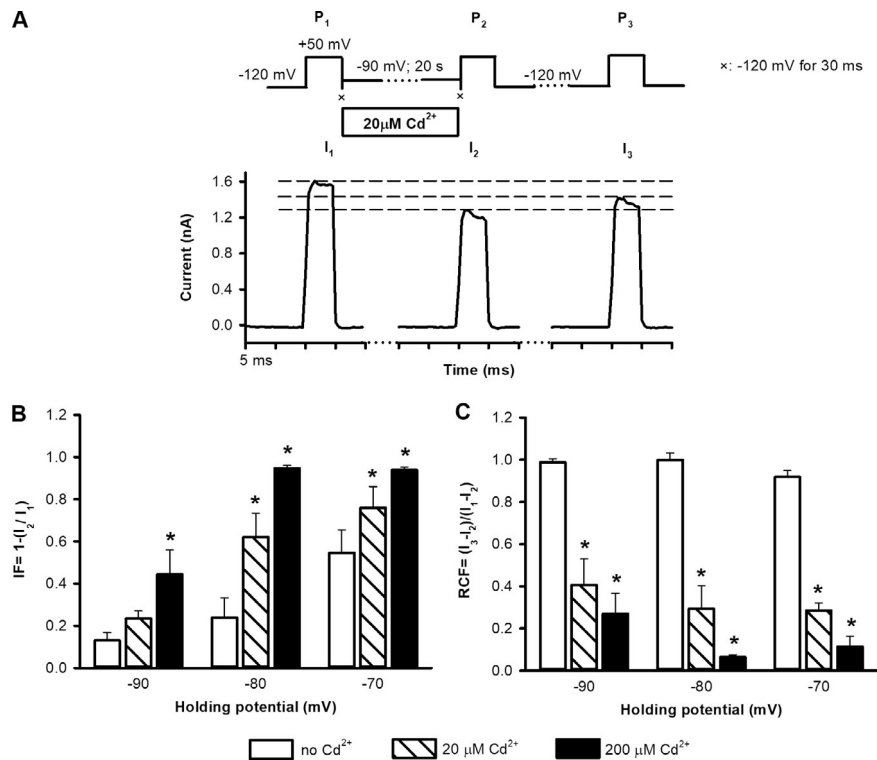


Figure 5. Determining the IFs and RCFs with intracellularly applied Cd²⁺. (A) An inside-out patch expressing T449A/V474C channels was exposed to 20 μM Cd²⁺ between P₁ and P₂ at -90 mV holding potential (top; see details also in Fig. 3). The holding potential before P₃ was -120 mV. Symbol × indicates a 30-ms-long pulse to -120 mV. Dashed reference lines are drawn to I₁, I₂, and I₃. (B) Based on the peak currents for each pulse, inactivated current fractions (IF = 1 - [I₂/I₁]) were calculated for the indicated holding potentials in the absence of Cd²⁺ (empty bars) or in the presence of 20 μM Cd²⁺ (hatched bars) or 200 μM Cd²⁺ (solid bars). (C) RCFs (RCF = [I₃ - I₂]/[I₁ - I₂]) were determined for the indicated holding potentials between P₁ and P₂ in the absence of Cd²⁺ exposure (empty bars) or in the presence of 20 μM Cd²⁺ (hatched bars) or 200 μM Cd²⁺ (solid bars) during the holding between P₁ and P₂. Data are given as mean ± SEM (n = 5); *, significant differences when P < 0.05 throughout the figure.

C, hatched bars). Thus, pulse durations in the three-pulse protocol are properly designed to avoid current loss in the absence of Cd²⁺ or A-gate movement. Second, channels exposed intracellularly to 200 μM Cd²⁺ during the 20-s-long holding period at -120 mV also gave no reduction in current, since they remained in the closed state at -120 mV (Fig. 4 A); IF ~0 and RCF ~1 (Fig. 4, B and C, black bars). These data demonstrate that the timing of the start/stop of the Cd²⁺ application is optimized to avoid Cd²⁺ exposure of the open channels when they are closing (end of P₁) or when closed channels are opening (during P₂); Cd²⁺ exposure is thus well controlled during the pulse protocol. A third control shows that all channels that are inactivated at -90-mV holding potential between P₁ and P₂ (Fig. 4 B, empty bar; IF significantly increased; P < 0.05) recover from inactivation in the absence of Cd²⁺ (Fig. S1 B and Fig. 4 C, empty bar; RCF ~1).

To assess the dependence of RCF on the holding potential, we chose holding potentials of -90, -80, or -70 mV between P₁ and P₂ in the presence of 20 or 200 μM intracellular Cd²⁺. Representative current traces of an experiment using 20 μM Cd²⁺ and -90-mV holding potential are shown in Fig. 5 A (the same experiment using 200 μM Cd²⁺ is in Fig. S2). Although no macroscopic current can be detected at -90 mV (Fig. 2, A and B), a small degree of steady-state inactivation occurs (IF, Fig. 5 B). At -80 and -70 mV, current increases (Fig. 2 A, inset), and inactivation becomes progressively more prominent (Fig. 2 D). IF increases with depolarizing holding potentials (Fig. 5 B) both in the absence and presence of Cd²⁺. In contrast, RCF decreased significantly only if Cd²⁺ was present during the holding potential between P₁ and P₂ (Fig. 5 C). This indicates that A-gate openings occur in the voltage range of -90 to -70 mV. Consequently, Cd²⁺ accesses 474C despite the lack

of macroscopic current at -90 mV (Fig. 2 A, inset). The RCF is significantly decreased in the presence of 20 μM and 200 μM Cd²⁺ (Fig. 5 C) regardless of the holding potential. In the absence of Cd²⁺, RCF values are essentially the same at all holding potentials (Fig. 5 C).

To correlate the increase in the current loss with Cd²⁺ exposure, we applied 20 μM Cd²⁺ for short (200 ms) durations over many cycles at -90 or -80 mV (Fig. 6). The pulse sequence (Fig. 6, A and B) started with a 5-ms depolarization to +50 mV to obtain the peak current I₁, then the patch was held at -120 mV for 60 s. This was followed by the unit cycle, which consisted of 200 ms at -90 mV (or -80 mV) and 800 ms at -120 mV. The unit cycle was repeated 200, 100, or 50 times, followed by a 60-s holding at -120 mV in the absence of Cd²⁺, and the sequence was terminated by a 5-ms depolarization to +50 mV to obtain I₂. Cd²⁺ was either absent (Fig. 6 A) or present (20 μM intracellular Cd²⁺, Fig. 6 B) during the indicated number of unit cycles. The cumulative fraction of channels modified by Cd²⁺ was calculated as IF = 1 - (I₂/I₁), as above. As indicated in Fig. 6, C and D, modification of 474C by cumulative exposure to Cd²⁺ is time and voltage dependent. A step to -80 mV produces more modification and increased IF than -90 mV. The pulse sequence gave cumulative Cd²⁺ exposure at 100 cycles similar to that obtained with a single 20-s application of 20 μM Cd²⁺. In summary, our results show that at holding potentials at which substantial steady-state inactivation occurs, the A-gate opens despite the lack of observable currents. This finding indicates that the CI state is populated through the C→O→OI pathway instead of the direct C→CI transition (see Discussion), as suggested by others (Klemic et al., 2001; Claydon et al., 2007, 2008; Jamieson and Jones, 2014).

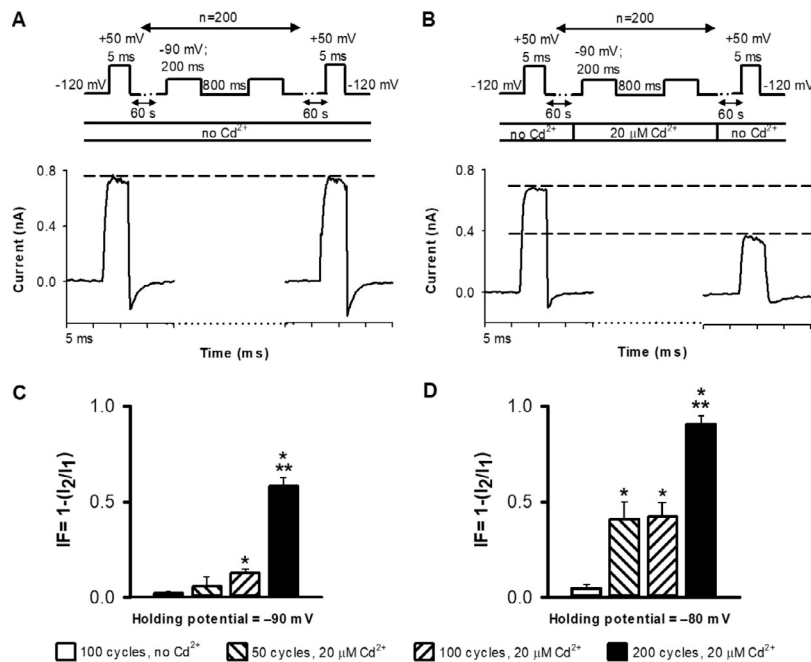


Figure 6. Cd²⁺ modification of T449A/V474C channels depends on the membrane potential and the duration of application. (A) Representative control pulse sequence for cumulative Cd²⁺ exposure. The unit cycle is defined as 200 ms at -90 and 800 ms at -120 mV holding potential. This cycle was repeated 200 times ($n = 200$). The cycles were flanked by 5-ms-long depolarizations to +50 mV in to obtain I_1 and I_2 , the peak inside-out patch currents before and after the cycles, respectively. 60-s-long holdings at -120 mV followed the first pulse to +50 mV and preceded the second pulse to +50 mV. Empty bar under the pulse protocol indicates that no Cd²⁺ was applied in the control sequence. Dashed reference line is drawn at I_1 . (B) Representative pulse sequence for cumulative Cd²⁺ exposure. The unit cycle was the same as in A except Cd²⁺ was applied as indicated by the bar (20 μM Cd²⁺). This cycle was repeated by $n = 200$ times. Dashed reference lines are drawn at I_1 and I_2 . (C) $1 - (I_2/I_1)$ at -90 mV holding (during the 200-ms segment of a unit cycle) at various cycle numbers in the absence of Cd²⁺ (empty bar, 50 cycles) or in the presence of 20 μM Cd²⁺ with 50 (left hatched bar), 100 (right hatched bar), and 200 cycles (filled bar). (D) Same experiment as in C except the holding potential was -80 mV during the 200-ms segment of a unit cycle. Mean values of $n = 3-7$ independent measurements are shown; error bars indicate SEM. * indicates that data are different from controls ($P < 0.05$). ** indicates that the data obtained with 200 cycles are also different from those obtained with 50 and 100 cycles ($P < 0.05$).

Locked-open A-gate prevents recovery from C-type inactivation

The experiments above suggest that an open A-gate is important for inactivation to occur at negative holding potentials. The ability of the channels to recover from inactivation may also mechanistically depend on the conformational status of the A-gate. Locking the A-gate in the open state is a suitable method to investigate determinants for recovery from inactivation.

Previous studies have shown that a *Shaker*-IR 476C mutant channel can be trapped in the open state by Cd²⁺ application from the intracellular side (Webster et al., 2004). This is caused by formation of a metal bridge between 476C in one subunit and a native histidine (H486) in a neighboring subunit, thus inhibiting A-gate closing even at negative voltages. A biophysical characterization of the T449A/V476C *Shaker*-IR channel expressed in HEK293 cells using inside-out patch-clamp configuration is given in Fig. S3. We designed control measurements to confirm the applicability of the protocol to study the consequence of A-gate locking by Cd²⁺ exposure on recovery from inactivation.

In T449A/V476C channels the extent of recovery from inactivation was monitored by current amplitudes evoked by 2-s-long repeated depolarizations to +50 mV from a holding potential of -120 mV. The ipi was set to 60 s (Fig. 7 A). Currents were normalized to the peak of the first pulse and plotted as a function of time elapsed from the first depolarizing pulse (Fig. 7 B). Repeated depolarizations in the absence of Cd²⁺ resulted in similar current amplitudes for the T449A/V476C channels; thus, the ipi is sufficient for full recovery of the current.

A second control experiment shown in Fig. 7, C and D, demonstrates the lack of peak current reduction by Cd²⁺ in the absence of 476C. The pulse protocol was identical to the one used

in Fig. 7 A except Cd²⁺ (20 μM) was applied to OI channels for 1 s. We recorded identical peak currents even after long (≤ 8 s) cumulative exposure of T449A/V476 channels to Cd²⁺ in the OI state (Fig. 7 D). Thus Cd²⁺ does not influence the currents in the absence of 476C.

To determine the dependence of recovery from inactivation on the status of the A-gate, we locked open the A-gate of T449A/V476C *Shaker*-IR channels by intracellular Cd²⁺ exposure (20 μM) after the full inactivation of the current (i.e., in the OI state) at +50 mV (Fig. 8, A and B). To do this, we essentially repeated the control experiment in Fig. 7 C except that T449A/V476C *Shaker*-IR channels were used. Cd²⁺ exposure was timed to the final 1,000 ms of the first 2-s-long depolarizing pulses (Fig. 8 A). The T449A/V476C *Shaker* construct inactivates completely by the time the Cd²⁺ application starts (Figs. 2 C and 8 A; $\tau_1 = 136$ ms; Cd²⁺ application starts after approximately $>7\times$ the time constant for inactivation). During the remaining time, patches were perfused by standard intracellular solution. This resulted in complete loss of peak current. Thus, channels that have A-gates locked open with Cd²⁺ are unable to recover from inactivation ($n = 5$). We suggest that these channels were trapped in the OI state and therefore remained nonconductive even after a 60-s-long holding at -120 mV. The loss of the current was irreversible. The kinetics of current loss was determined by repeated exposure of the OI channels to 20 μM Cd²⁺ for short, 200-ms durations (Fig. 8 B). This resulted in a quantized and significant decrement of peak currents (Fig. 8 C) upon cumulative exposure. The loss of the current as a function of cumulative modification time followed a single exponential decay, indicating a specific interaction between the cysteine residue at position 476C and Cd²⁺.

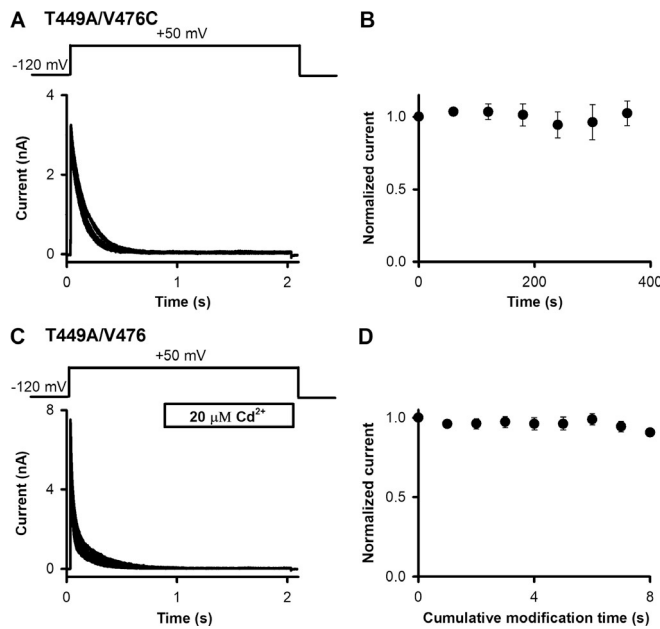


Figure 7. Control measurements for determining OI→O transition. (A) An inside-out patch pulled from a HEK293 cell expressing T449A/V476C *Shaker*-IR channels was repeatedly depolarized from a holding potential of -120 mV to $+50$ mV for 2 s using the pulse protocol shown above the superimposed raw current traces. The ipi was 60 s. (B) The peak current for each pulse was determined and normalized to the peak current of the first pulse, and then these were plotted as a function of the time elapsed from the first depolarizing pulse. (C) For T449A/V476 channels, the same protocol was used as in A, except that the intracellular surface of the patch was exposed to $20 \mu\text{M Cd}^{2+}$. The timing (applied to fully inactivated channels 1 s after the depolarization to $+50$ mV) and the duration (1 s) of the Cd^{2+} application are indicated by the horizontal bar. (D) Peak currents for each pulse were determined, normalized to the peak current of the first pulse, and plotted as a function of the cumulative exposure time to $20 \mu\text{M Cd}^{2+}$. Error bars indicate SEM throughout the figure; $n = 5$.

To confirm that this nonconductive OI state is created by a Cd^{2+} bridge between 476C and H486 in our inactivating V449A construct, we tested whether the Cd^{2+} -induced current loss could be prevented by the protonation of H486 at low pH. A deprotonated H486 is critical for the metal-bridge formation by Cd^{2+} with 476C (Holmgren et al., 1998; Webster et al., 2004). Depolarizing pulses to $+50$ mV were applied from a holding of -120 mV every 15 s. Cd^{2+} application started 200 ms before each pulse and ended simultaneously with the depolarization. The pH of Cd^{2+} containing intracellular solutions was set to 7.36 (Fig. 9 A) or 5.3 (Fig. 9 B). The development of the locked-open state at pH 7.36 (Fig. 9 A) is indicated by the increased relative amplitudes of the tail currents (compared with tail currents at pH 5.3), followed by a slow decay corresponding to inactivation. The trace-to-trace decrease in the peak currents at $+50$ mV indicates that locked-open and inactivated channels do not recover from inactivation (Fig. 9 A). This current phenotype and the inactivation process of open channels with locked-open A-gates have been described previously for V476C mutant *Shaker* channels (Liu et al., 1997; Holmgren et al., 1998; Webster et al., 2004). In contrast, at pH 5.3, where the metal bridge formation is inhibited, both the shape of the tail currents and peak amplitudes

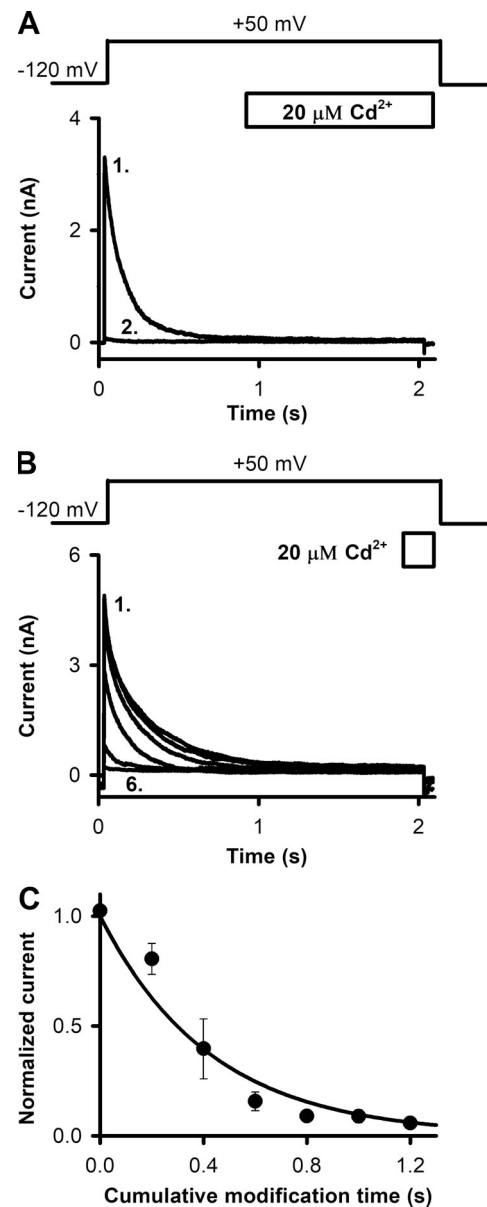


Figure 8. Determining the extent of recovery from inactivation after locking open the activation gate of T449A/V476C *Shaker*-IR channels by Cd^{2+} . (A and B) Inside-out patches pulled from HEK293 cells expressing T449A/V476C *Shaker*-IR channels were repeatedly depolarized from a holding potential of -120 mV to $+50$ mV for 2 s using the pulse protocols shown above the corresponding raw current traces. The ipi was 60 s. The timing and duration of the exposure of the intracellular surface of the patch to $20 \mu\text{M Cd}^{2+}$ are indicated by the empty bars. (A) 1,000-ms-long perfusion with $20 \mu\text{M Cd}^{2+}$ started 1,000 ms after the start of each depolarizing pulse, and the patch was perfused with standard intracellular solution for the remaining time and during the interpulse holding potential. Numbers in the panel indicate the first and second traces. (B) 200-ms-long perfusion with $20 \mu\text{M Cd}^{2+}$ started 1,800 ms after the start of each depolarizing pulse, and the patch was perfused with standard intracellular solution for the remaining time and during the interpulse holding potential. Numbers in the panel indicate the first and last traces. (C) Peak currents were determined from experiments as in B and were normalized to the peak of the first pulse (in the absence of Cd^{2+}) and plotted as a function of cumulative exposure to $20 \mu\text{M Cd}^{2+}$. Solid line indicates the best-fit single exponential function to the averaged data points. Mean \pm SEM values for $n = 5$ experiments are shown.

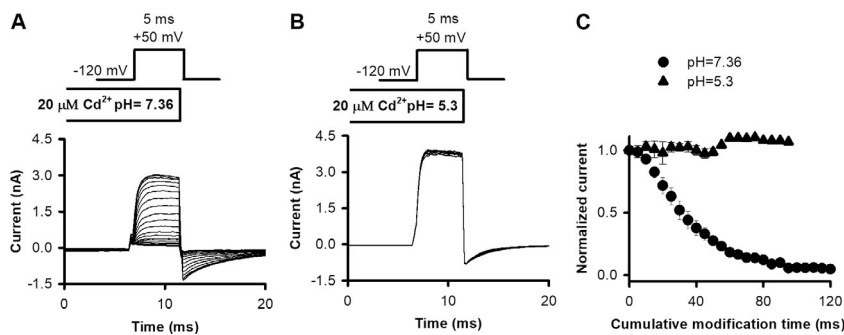


Figure 9. Applying Cd²⁺ in low-pH solutions inhibits metal bridge formation. (A and B) Inside-out patches expressing T449A/V476C channels were repeatedly depolarized the holding potential of -120 mV to +50 mV every 15 s in the presence of 20 μM Cd²⁺ at pH 7.36 (A) and pH 5.3 (B). Cd²⁺ was applied intracellularly for 200 ms before opening and during the depolarization to +50 mV. Pulse protocols, with appropriate ipi, were run three to four times in the absence of Cd²⁺ to verify the stability of the peak currents (not depicted). (C) Peak currents for each pulse were normalized to the peak current of the first pulse and plotted as a function of cumulative modification time. Results represented by different symbols were obtained at pH 7.36 (solid circles) and pH 5.3 (solid triangles). Error bars indicate SEM for $n = 4-5$ independent measurements.

of the currents at +50 mV are conserved during repeated depolarizations (Fig. 9 B). Normalized peak currents decrease during cumulative application of Cd²⁺ at pH 7.36 (solid circles); however, at pH 5.3 (solid triangles), the peaks were maintained at a steady level (Fig. 9 C; $n = 4-5$).

Discussion

The thermodynamic four-state cycle (Fig. 1) is a useful approximation for exploring gating mechanisms. This model obtains across all voltages, but the probabilities of state occupancy and kinetics of gating transitions will be voltage dependent and thereby determine the likely pathways taken to a given state. In this study, we explored electrophysiologically “silent” gating transitions (transitions not accompanied by measurable current) of *Shaker*-IR during the development of steady-state inactivation and recovery from inactivation. In both inactivation and recovery from inactivation, entry into and exit from the CI state is crucial. At negative membrane potentials, we find that *Shaker* K⁺ channels populate the CI state through the C→O→OI→CI pathway and that the OI→CI transition is mandatory before recovery from inactivation. Support for these conclusions was achieved by demonstrating that (1) the A-gate opens at membrane potentials at which the CI state develops and (2) recovery from inactivation can be prevented when the A-gate is locked open in inactivated channels (OI-state stabilized). Specifically, in the first case, we found that Cd²⁺ modifies the T449A/V474C *Shaker*-IR channels at -90 mV, a membrane potential at which steady-state inactivation develops in the absence of measurable current, i.e., the opening of the A-gate of the channel allows Cd²⁺ to reach 474C behind the A-gate. In the second case, locking the A-gate open using a Cd²⁺ bridge between 476C and H486 in the T449A/V476C construct renders the channels permanently inactivated, thereby impeding recovery from inactivation. Current literature shows a strong allosteric coupling between the activation and inactivation gates of *Shaker*: opening of the A-gate promotes closure of the I-gate at the selectivity filter. Here we show that an open A-gate can lock the I-gate in the nonconducting conformation as well. Considering the facilitation of I-gate closure by an open A-gate, we propose that the CI state is populated through the C→O→OI→CI pathway in *Shaker* at negative membrane potentials.

Closed-state inactivation has been described and reviewed for many ion channels, including voltage-gated sodium, calcium, and potassium channels (Bähring and Covarrubias, 2011; Bähring et al., 2012). The existence of closed-state inactivation in *Shaker*-IR channels has been shown in the ILT mutant at low pH, but not at physiological pH (Claydon et al., 2007, 2008). The ILT triple substitution in S4, however, disrupts fundamentally the cooperativity between the voltage-sensor and the pore domain and may affect the communication between the activation and C-type inactivation gates as well (Smith-Maxwell et al., 1998). A multiple-state inactivation model was described for the L382I mutant of *Shaker* which, in contrast to *Shaker*-IR channels, can inactivate directly from one of the four deep closed states (Ayer and Sigworth, 1997). Residue 382 is located in the S4-S5 linker close to the bottom of the S4 helix and has been identified as an important element in the coupling between activation and inactivation gates (Long et al., 2005); thus the mutation can generate new possible transitions between the C and CI states. Closed-state inactivation has been proposed for *Shaker*-IR T442A channels also, which contain a T-to-A mutation in the second threonine residue in TTVGYGD signature sequence (Labro et al., 2018). This mutation largely affects the gating properties of the channel because the A-gate opening does not lead to C-type inactivation (inverted allosteric coupling between the activation and C-type inactivation gates; Labro et al., 2018).

In contrast to experiments using mutant *Shaker* channels having dramatically altered the properties mentioned above (ILT, L382I, and T442A), which led to conclusions based on ionic current measurements at small driving force, we improved the experimental settings in two areas. First, we used *Shaker* T449A/V474C and T449A/V476C mutants, in which the natural gating processes are conserved, and thus the gating parameters are comparable to the parameters of the wild-type channel (Fig. 2 and Fig. S3). Second, we monitored and manipulated directly the status of the A-gate by Cd²⁺ modification. This method provides a delicate tool to track A-gate opening (Liu et al., 1997) and lock it in the open state (Holmgren et al., 1998; Webster et al., 2004). Cysteine modification and metal-bridge formation by Cd²⁺ for 474C and 476C channels are based on sensitive and specific chemical interactions. Because of the applied experimental strategies, the electrophysiologically silent transitions were revealed,

and the fast perfusion system provided a precise way to perform Cd^{2+} modification assays.

In their seminal paper, [del Camino and Yellen \(2001\)](#) showed that the Cd^{2+} modification rate of *Shaker* 474C is in a perfect agreement with the normalized open probability (P_o) of the channel, i.e., the A-gate provides gated access for Cd^{2+} to the cysteine in position 474C, located behind the A-gate. As the P_o -V relationship is a continuous function, we aimed at increasing the reliability of the RCF readout ([Fig. 5](#)) by an appropriate negative control. We showed that when a holding potential of -120 mV was used between P_1 and P_2 , very long exposure (20 s) of high concentration (200 μM) of Cd^{2+} (10 \times the concentration used in [Fig. 5 A](#)) resulted in $\text{RCF} = 1$, i.e., the data shown in [Fig. 4 C](#) serve as a good negative control for the data in [Fig. 5](#). Although *Shaker*-IR channels in the closed state have a miniature, but measurable, conductance (upper limit of 16 fS; [Soler-Llavina et al., 2003](#)), our negative control in [Fig. 4](#) argues that this conductance does not allow significant modification of 474C by Cd^{2+} at -120 mV.

Based on the experiments shown in [Figs. 4 and 5](#) and the analysis above, we attribute the reduction of RCF to the interaction of Cd^{2+} with cysteines behind the A-gate. The specific interaction between Cd^{2+} and cysteine is supported by the fact that a higher concentration of Cd^{2+} resulted in larger reduction of RCF, as expected for a pseudo-first-order chemical reaction between cysteine and Cd^{2+} ([Fig. 5](#)). In addition, repeated brief changes in the holding potential to -90 mV (or -80 mV; [Fig. 6, C and D](#)) in the presence of Cd^{2+} resulted in a similar increase in IF as a single, long-duration application of Cd^{2+} ([Fig. 5](#)). Furthermore, the analysis of the time dependence of the reaction could be extended by repeated short-duration application of Cd^{2+} without increasing the chance of nonspecific interactions, which may happen at a single, long-lasting application of Cd^{2+} . In summary, we conclude from experiments shown in [Figs. 3, 4, 5, and 6](#) that the Cd^{2+} accessibility assay based on a specific interaction between 474C and Cd^{2+} and performed by a fast-step perfusion system in our experimental setup provides an appropriate tool to detect A-gate opening via determining the reduction in RCF.

Data in [Fig. 5](#) show that RCF is nonzero at all tested membrane potentials. This may mean that either only a fraction of channels that contribute to steady-state inactivation are in the state where Cd^{2+} can interact with them, or perhaps all of them are in the Cd^{2+} -accessible state but Cd^{2+} does not modify all the channels within 20 s. In the latter case, the RCF should decrease with an increase in the Cd^{2+} concentration, as we have shown clearly in [Fig. 5 C](#). Short, repeated applications of Cd^{2+} resulting in increased cumulative exposure time resulted in more significant current loss (increase in IF, [Fig. 6](#)), which also supports the hypothesis that channels are in the Cd^{2+} -accessible state when steady-state inactivation takes place. Both IF and RCF parameters are expected to show membrane potential dependence, which is demonstrated in [Figs. 5 and 6](#).

Our experiments clearly demonstrate that a significant fraction of *Shaker*-IR channels have open A-gates and interact with Cd^{2+} at negative membrane potentials where steady-state inactivation takes place. A simple conclusion from this finding could be that some channels open and then inactivate at -90 mV,

and these are the channels that can be modified by Cd^{2+} . That is, opening of the channels precedes inactivation, and the gating steps resulting in steady-state inactivation are $\text{C} \rightarrow \text{O} \rightarrow \text{OI}$. In reality, both O and OI channels can be modified by Cd^{2+} through 474C ([Panyi and Deutsch, 2007](#)). Thus, we cannot rule out the possibility that the reduction of RCF is the consequence of Cd^{2+} interacting with channels generated by $\text{C} \rightarrow \text{CI} \rightarrow \text{OI}$ transition at -90 mV. However, we suggest this latter scenario is less likely for two reasons. First, Cd^{2+} modification of OI channels occurs at an ~ 20 -fold lower rate compared with O channels ([Panyi and Deutsch, 2007](#)). Second, recent publications argue that opening of the A-gate triggers the collapse of the selectivity filter in *Shaker*. That is, the principal transition leading to the OI state is $\text{O} \rightarrow \text{OI}$ ([Labro et al., 2018](#)), so the state that was modified by Cd^{2+} should rather be the O state preceding the OI state. This conclusion may hold for *Shaker* constructs, e.g., T449A, where C-type inactivation dominates the inactivation process ([Jamieson and Jones, 2014](#)). However, in *Shaker*-IR channels displaying other slow-inactivation mechanisms, e.g., U-type inactivation, transitions to the inactivated state may occur predominantly from closed states ([Klemic et al., 2001](#)).

Compared to other gating transitions, relatively less is known about the pathway and the molecular rearrangements leading to recovery from C-type inactivation. In a classic scenario during prolonged depolarization, the slow inactivation gate shuts (OI state), and returning to a negative membrane potential induces the $\text{OI} \rightarrow \text{CI}$ transition. Channels in the CI-state will, under a sufficiently negative and prolonged repolarization period, recover from inactivation and populate the closed state (C). The rate of recovery can be modulated by many factors, e.g., the concentration of extracellular and intracellular cations, extracellular pH, and the membrane potential ([Pardo et al., 1992](#); [Levy and Deutsch, 1996a, 1996b](#); [Rasmusson et al., 1998](#); [Ray and Deutsch, 2006](#)). The rate-limiting step in recovery from inactivation is the $\text{CI} \rightarrow \text{C}$ transition, as $\text{OI} \rightarrow \text{CI}$ is very fast (milliseconds) in *Shaker* ([Panyi and Deutsch, 2006](#)), whereas recovery from inactivation is on the time scale of many seconds. Our previous finding that the A-gate of inactivated channels closes relatively quickly (time constant is 23 ms at -120 mV; [Panyi and Deutsch, 2006](#)) does not rule out the possibility of recovery (i.e., a selectivity filter in the conductive state) with an open A-gate.

If closure of the A-gate is obligatory for recovery, then channels with a locked-open A-gate will remain inactivated, i.e., recovery from inactivation is precluded. This is what we demonstrated for *Shaker*-IR in [Fig. 8](#): locking the A-gate open in inactivated channels (OI) using a single, 1-s application of Cd^{2+} or by repeated administration for 200-ms-long Cd^{2+} pulses completely prevented recovery from inactivation. The decrement in peaks upon cumulative Cd^{2+} exposure follows a single-exponential decay, indicating the specific interaction between Cd^{2+} and T449A/V476C channels in [Fig. 8 C](#). We concluded from these data that recovery from inactivation is not possible with a locked-open A-gate.

This conclusion is based on the loss of the current upon Cd^{2+} application ([Fig. 8](#)). Divalent cations are known to inhibit ion channels, e.g., Ba^{2+} blocks *Shaker* and other K_v channels

(Armstrong and Taylor, 1980; Hurst et al., 1995; Harris et al., 1998; Gibor et al., 2004), and Cd^{2+} blocks several ion channels (Hobai et al., 1997; Abbruzzese et al., 2010; Tu et al., 2016). A critical control to support our conclusion was therefore to demonstrate that Cd^{2+} does not inhibit *Shaker*-IR lacking the critical cysteine in position 476 (Fig. 7 B). We also showed that the 60-s-long ipi between repetitive depolarizing pulses is sufficient for a full recovery from inactivation in control (Fig. 7 A). Thus, the current loss demonstrated in Fig. 8 cannot be explained by an inadequately short holding at -120 mV between pulses.

Locking the A-gate in the open configuration using a Cd^{2+} bridge between 476C and a histidine of a neighboring subunit (H486) was described in detail for a *Shaker*-IR channel (Webster et al., 2004). To confirm the existence of the Cd^{2+} bridge between H486 and 476C in the inactivating T449A/V476C *Shaker*-IR construct, we changed the intracellular pH. At low intracellular pH (5.3), where histidines are protonated ($\text{pK}_a \sim 6.0$), Cd^{2+} is unable to modify the currents (Fig. 9 C). This argues for a specific action between Cd^{2+} , H486, and 476C.

We conclude that closure of the A-gate is a precondition for recovery from inactivation and that the sequence of gating transitions is $\text{O} \rightarrow \text{OI} \rightarrow \text{CI} \rightarrow \text{C}$ at negative membrane potentials. Based on this scheme, an intracellular open channel blocker that prevents A-gate closure might slow and delay recovery from inactivation. Analogous to our findings in *Shaker*, it has been demonstrated for Na_v channels that channel deactivation needs to precede recovery from inactivation (Kuo and Bean, 1994).

Our investigation has used a novel combination of voltage protocols and modification strategies to define the gating scheme that underlies steady-state inactivation for *Shaker* at negative membrane potentials, and to identify a prerequisite for recovery from inactivation. We propose that the most likely pathway for the development of steady-state inactivation is $\text{C} \rightarrow \text{O} \rightarrow \text{OI} \rightleftharpoons \text{CI}$. Moreover, the $\text{OI} \rightarrow \text{CI}$ transition is a prerequisite for recovery from inactivation. This sequence of gating transitions complies with tight coupling between the activation and C-type inactivation gates. The biological significance of these findings is that both steady-state inactivation and recovery from inactivation essentially determine the fraction of channels available to open, which is critical to excitability in nerve and muscle cells.

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Supplemental material

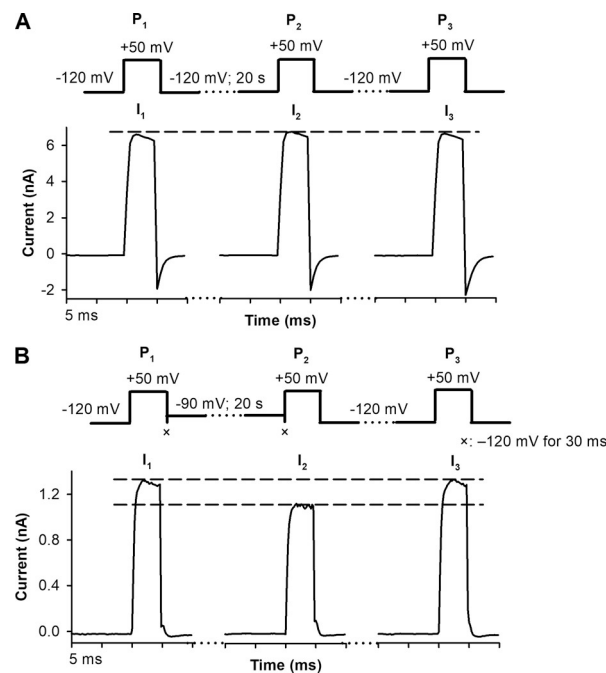


Figure S1. **Control experiments to validate the suitability of the three-pulse protocol.** Inside-out patch currents were recorded in HEK293 cells expressing T449A/V474C *Shaker*-IR channels. **(A)** Representative traces show currents evoked by three 5-ms-long depolarizing pulses to +50 mV from the holding potential of -120 mV. The IF and RCF values calculated from I_1 , I_2 , and I_3 peak currents are in Fig. 4, B and C. **(B)** Representative traces show currents evoked by three 5-ms-long depolarizing pulses to +50 mV. The holding potential was -120 mV before P_1 and P_3 . The holding potential between pulses P_1 and P_2 was -90 mV. The IF and RCF values calculated from I_1 , I_2 , and I_3 peak currents are in Fig. 4, B and C.

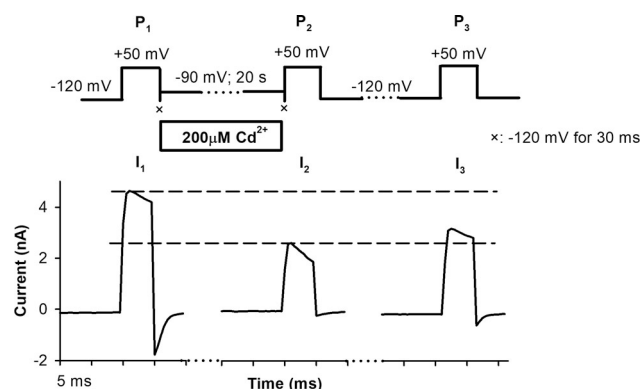


Figure S2. **Exposure of an inside-out patch to 200 μM Cd^{2+} at -90 mV in a three-pulse protocol.** Inside-out patch currents were recorded in HEK293 cells expressing T449A/V474C *Shaker*-IR channels. Representative traces show currents evoked by three depolarizing pulses to +50 mV. The holding potential was -120 mV before P_1 and P_3 . The holding potential between pulses P_1 and P_2 was -90 mV. 200 μM Cd^{2+} was applied intracellularly (empty horizontal bar). The IF and RCF values calculated from I_1 , I_2 , and I_3 peak currents are in Fig. 5, B and C.

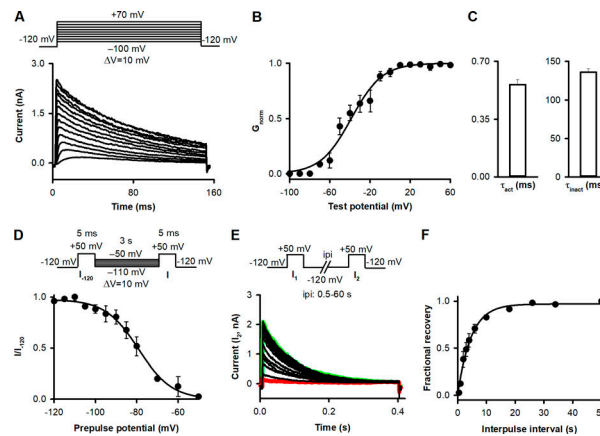


Figure S3. Biophysical characterization of T449A/V476C Shaker-IR ion channels. (A) To characterize the steady-state activation of T449A/V476C mutant channels, inside-out patches were held at -120 mV and depolarized to test potentials ranging from -100 to $+70$ mV in steps of 10 mV every 60 s. The duration of the depolarizing pulses was 100 ms (see protocol above). Representative current trace is shown in A. (B) G - V data were obtained from the peak currents (I_{peak}) at test potential (V_{test}) and the K^+ reversal potential (E_{rev}) using $G = I_{\text{peak}} / (V_{\text{test}} - E_{\text{rev}})$. Then G values were normalized for the maximum conductance (G_{norm}) and plotted as a function of test potential. The superimposed solid line shows the best-fit Boltzmann function to the mean values \pm SEM ($n = 5$). The Boltzmann function was fitted to the G - V relationship obtained from $n = 5$ patches to result in $V_{1/2} = -38.8 \pm 2.1$ mV, $k = 15.3 \pm 1.8$ mV ($n = 5$, mean \pm SEM). (C) To determine the activation time constant (τ_{act}) we applied 5 -ms-long depolarizing pulse from a holding potential of -120 mV to $+50$ mV. Current traces were fitted using the Hodgkin-Huxley n^4 -model; τ_{act} was used to describe the activation kinetics. Inactivation time constant (τ_{inact}) of the current at $+50$ mV was determined by fitting a single exponential function to the decaying part of the currents evoked by a $2,000$ -ms-long depolarizing pulse. (D) To describe the voltage dependence of steady-state inactivation from a holding potential of -120 mV, the patch was stepped to prepulse potentials between -110 and -50 mV ($\Delta V = 5$ mV) for 3 s, then a 5 -ms-long test pulse to $+50$ mV was applied to elicit K^+ currents (see protocol above). The fraction of noninactivated channels at each voltage was calculated as I/I_{-120} , where I is the peak current evoked by 5 -ms depolarization to $+50$ mV from a prepulse potential applied for 3 s, and I_{-120} is the peak current evoked by identical depolarization from the holding potential of -120 mV. The superimposed solid line shows the best-fit Boltzmann function to the mean values of $I/I_{-120} \pm$ SEM ($n = 5$). The Boltzmann function was fitted to the voltage dependence of steady-state inactivation relationships of $n = 5$ patches to result in $V_{1/2} = -79.1 \pm 0.7$ mV, $k = -8.0 \pm 0.6$ mV ($n = 5$, mean \pm SEM). (E) For measuring the kinetics of recovery from inactivation, pairs of depolarizing pulses were delivered from the holding potential of -120 to $+50$ mV for 200 ms (see protocol above). Currents recorded during the second pulses (I_2) are shown at various ipi at -120 mV holding. Traces recorded with ipi of 0.5 and 60 s are highlighted in red and green, respectively. (F) The fractional recovery was calculated as $(I_2 - I_{\text{SS1}}) / (I_1 - I_{\text{SS1}})$, where I_2 and I_1 are the peak currents during the second and first pulse, respectively, and I_{SS1} is the steady-state current at the end of the first depolarization. Data points were fitted with by single exponential function to give the time constant of recovery from inactivation (τ_{rec}). Data are given as mean \pm SEM; $\tau_{\text{rec}} = 4.6 \pm 0.6$ s ($n = 3$).