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Extracting the speed of sound in quark–gluon plasma with ultrarelativistic lead–lead collisions at the LHC

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Abstract

Ultrarelativistic nuclear collisions create a strongly interacting state of hot and dense quark–gluon matter that exhibits a remarkable collective flow behavior with minimal viscous dissipation. To gain deeper insights into its intrinsic nature and fundamental degrees of freedom, we determine the speed of sound in an extended volume of quark–gluon plasma using lead–lead (PbPb) collisions at a center-of-mass energy per nucleon pair of 5.02 TeV. The data were recorded by the CMS experiment at the CERN LHC and correspond to an integrated luminosity of 0.607 nb^{-1}. The measurement is performed by studying the multiplicity dependence of the average transverse momentum of charged particles emitted in head-on PbPb collisions. Our findings reveal that the speed of sound in this matter is nearly half the speed of light, with a squared value of $0.241 \pm 0.002\,(\mathrm{stat}) \pm 0.016\,(\mathrm{syst})$ in natural units. The effective medium temperature, estimated using the mean transverse momentum, is $219 \pm 8\,(\mathrm{syst})\,\mathrm{MeV}$. The measured squared speed of sound at this temperature aligns precisely with predictions from lattice quantum chromodynamic (QCD) calculations. This result provides a stringent constraint on the equation of state of the created medium and direct evidence for a deconfined QCD phase being attained in relativistic nuclear collisions.

Keywords: CMS, quark–gluon plasma, speed of sound, ultra-central, QCD equation of state

1. Introduction

When heavy atomic nuclei collide at relativistic speeds, a transformation occurs, giving rise to an exotic state of matter with a temperature above several trillion kelvin and known as the quark–gluon plasma (QGP) [1–4]. In this realm of extreme temperatures, quarks and gluons break free from their confined existence inside hadrons, traversing long distances (e.g. several fm) compared to the size of individual nucleons. The emergence of the QGP represents a fundamental prediction of quantum chromodynamics (QCDs) [5, 6], the theory that elucidates the nature of the strong force. More remarkably, this strongly interacting QGP matter is found to exhibit the characteristics of an almost 'perfect liquid' with little frictional momentum dissipation [7–10]. Its collective dynamics and macroscopic properties are well described by the principles of nearly ideal relativistic hydrodynamics.

The equation of state (EoS) reveals the underlying fundamental degrees of freedom of a substance and is an invaluable tool to infer how the substance will respond to changes in its energy density. In fluid-like environments, the study of sound modes arising from longitudinal compression provides a means to determine the corresponding speed of sound, denoted as c_s. This parameter, whose square is defined as the rate of pressure P change in response to variations in energy density ε, $c_s^2 = \mathrm{d}P/\mathrm{d}\varepsilon$ [11], plays a pivotal role in characterizing the nature of the medium under investigation and in constraining models of corresponding EoS. The exploration of the sound wave propagation in strongly correlated systems, ranging from neutron stars to ultracold atomic gases [12, 13], has garnered significant interest in recent years. Various methodologies have been proposed to experimentally extract the speed of sound in a QGP fluid [14–18], offering a direct means to constrain the QCD EoS. Notably, constraints on the speed of sound in hot QCD matter have been inferred through a comparison of relativistic nuclear collision data with theoretical models within a Bayesian framework [15]. Recently, an effort to directly extract c_s^2 in the QGP phase was made by establishing a connection to an effective static, uniform fluid system [16]. That work was based on only two independent measurements of the charged-particle multiplicity density and mean transverse momentum (p_{T}) in lead–lead (PbPb) collision data from the ALICE experiment at center-of-mass energies per nucleon pair $\sqrt{s_{\mathrm{NN}}} = 2.76$ and $5.02\,\mathrm{TeV}$, and yielded a value of $c_s^2 = 0.24 \pm 0.04$ in natural units at a temperature of $222 \pm 9\,\mathrm{MeV}$. This result is in line with lattice QCD predictions, albeit subject to significant experimental uncertainties.

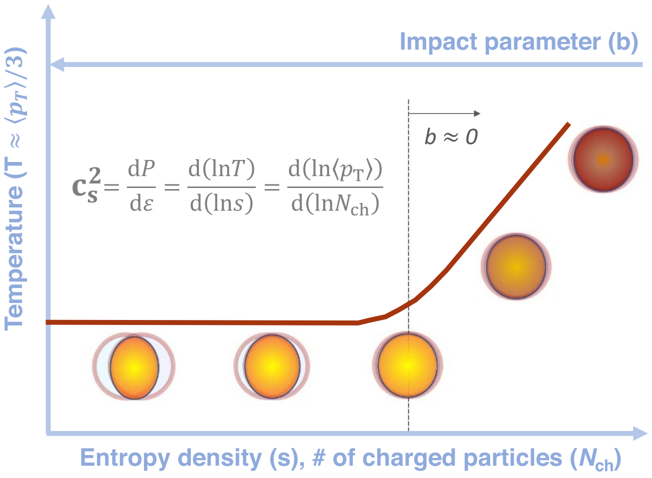

Figure 1. Conceptual representation of temperature vs. entropy density from mid-central to ultra-central heavy ion collisions.

To increase the precision by which the speed of sound can be determined, a new hydrodynamic probe was later proposed in [17] utilizing the multiplicity dependence of mean $\mathrm{p_T}$ measurements at a fixed $\sqrt{s_{\mathrm{NN}}}$. This innovative technique makes use of 'ultra-central' collisions in which the ions overlap almost entirely, i.e. collide at a very small impact parameter (b). A conceptual representation of this probe is illustrated in figure 1. The impact parameter of a heavy ion collision determines the size of the nuclear overlap region (system size), which is strongly correlated with the energy and entropy deposited in the initial state and the number of emitted charged particles in the final state ('multiplicity', N_{ch}). As the impact parameter decreases and collisions become increasingly central, both the system size and deposited energy increase, while maintaining a nearly constant initial energy density and temperature. However, this trend reaches its limit when $b \to 0$. In this case, the initial system size is limited by the sizes of the participating nuclei. For symmetric PbPb collisions, this would be the size of a Pb nucleus. More energy and entropy can still be deposited into the fixed volume through fluctuations in the number of interacting partons. By examining the response of the temperature T to the increasing entropy density s at $b \sim 0$, the speed of sound can be extracted based on fundamental thermodynamic laws,

$$c_s^2 = \frac{\mathrm{d}P}{\mathrm{d}\varepsilon} = \frac{s\mathrm{d}T}{T\mathrm{d}s} = \frac{\mathrm{d}\langle p_{\mathrm{T}}\rangle/\langle p_{\mathrm{T}}\rangle}{\mathrm{d}N_{\mathrm{ch}}/N_{\mathrm{ch}}}. \qquad (1)$$

Here, in terms of experimental observables, s is directly proportional to N_{ch}, while the temperature T relates to the average transverse momentum ($\langle p_{\mathrm{T}}\rangle$) of emitted particles with respect to the beam axis [16]. Full hydrodynamic simulations, such as those made possible using the TRAJECTUM model [19], have verified the above relationship, although there are features that are not captured, as will be discussed later. As the c_s^2 value depends only on the relative variation in $\langle p_{\mathrm{T}}\rangle$ and N_{ch}, any global changes to the observables, such as an increase in the system entropy through hadronic resonance decays [20], will not affect the result.

In this paper, we present a precise determination of the speed of sound in QGP using ultra-central PbPb collision data at $\sqrt{s_{\mathrm{NN}}} = 5.02\,\mathrm{TeV}$, collected in 2018 by the CMS experiment at the CERN LHC. By achieving a level of precision of several percent, comparable to theoretical uncertainties, our results serve as a robust benchmark for comparison with hydrodynamic simulations and lattice QCD calculations of the EoS. These comparisons provide the most stringent and direct constraints on the degrees of freedom attained by the medium created in these collisions. Tabulated results are provided in the HEPData record for this analysis [21].

2. The CMS detector

The CMS apparatus [22] is a multipurpose, nearly hermetic detector, designed to trigger on [23, 24] and identify electrons, muons, photons, and hadrons [25–27]. The initial triggering is done with the level-1 system, which uses customized hardware to make the rapid online decision whether or not to accept an event and deliver it to the second system, the high level trigger (HLT). The HLT uses a large CPU farm to perform optimized online event reconstruction and characterize an event. A global 'particle-flow' algorithm [28] aims to reconstruct all individual particles in an event, combining information provided by the all-silicon pixel and strip tracker, and by the crystal electromagnetic and brass-scintillator hadron calorimeters, operating inside a 3.8 T superconducting solenoid, with data from the gas-ionization muon detectors embedded in the flux-return yoke outside the solenoid. Hadron forward (HF) calorimeters [29], made of steel and quartz fibers, extend the pseudorapidity ($\eta = -\ln(\tan(\theta/2))$, where the polar angle θ is defined relative to the counterclockwise beam) coverage provided by the barrel and endcap detectors. Two zero-degree calorimeters (ZDCs) [30], made of quartz-fibers and plates embedded in tungsten absorbers, are used to detect neutrons from nuclear dissociation events.

3. Data samples, event reconstruction and selection

The data analyzed, before applying the selection described below, consist of 4.27 $\times 10^9$ minimum bias events, corresponding to an integrated luminosity of 0.607 nb^{-1}. The minimum bias events are triggered by requiring total energy signals above readout thresholds, which are in the range 6–12 GeV, on both sides of the HF calorimeters [24]. Beam-gas interactions and nonhadronic collisions are rejected by requiring the shapes of the clusters in the pixel tracker to be compatible with those expected from particles produced by a PbPb collision [31]. The events are also required to have at least one reconstructed primary vertex associated with two or more tracks within a distance of 15 cm from the nominal interaction point along the beam axis. The primary vertex is selected as the one with the highest track multiplicity in the event. Events with concurrent interactions per bunch crossing contribute to about 0.5% of the full data sample and are rejected based on the correlation of total energy deposited in the HF and ZDC detectors, following the procedure used in [32]. The collision centrality in PbPb events, i.e. the degree of overlap or impact parameter of the two colliding nuclei, is commonly determined by the total transverse energy deposit in both HF calorimeters, $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ [31]. As the main focus of this work is on collisions at small impact parameters, we analyzed only the 10% of PbPb events that had the largest $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$. This class contains the ultra-central collision events of interest.

To ease the computational load for high-multiplicity central PbPb collisions, track reconstruction for PbPb events is done in two iterations. The first iteration reconstructs tracks from signals ('hits') in the silicon pixel and strip tracker that are compatible with trajectories of particles with $p_{\mathrm{T}} > 1.0\,\mathrm{GeV}$, while the second iteration reconstructs tracks compatible with trajectories of particles with $0.3 < p_{\mathrm{T}} < 1.0\,\mathrm{GeV}$ using solely the pixel detector. In the analysis, the tracks have the additional selection requirement of $|\eta| < 0.5$ for the best tracking performance. More details on the track reconstruction and selection can be found in [33]. The tracking efficiency ($\varepsilon_{\mathrm{eff}}$) and misreconstruction rate ($\varepsilon_{\mathrm{mis}}$) are evaluated using the HYDJET [34] event generator, together with a full GEANT4 [35] simulation of the CMS detector response. These factors are combined to obtain an overall correction factor, $\varepsilon_{\mathrm{trk}} = \varepsilon_{\mathrm{eff}}/(1-\varepsilon_{\mathrm{mis}})$, which is used to account for detector effects on the total number of reconstructed tracks. The $\varepsilon_{\mathrm{trk}}$ factor is calibrated not only in terms of p_{T} and η, but also as a function of the detector occupancy. The occupancy is estimated by the total number of clusters registered in the silicon pixel tracker N_{pixel}, where a weak linear decline of $\varepsilon_{\mathrm{trk}}$ by up to 7% over an increase of N_{pixel} by 30% is observed. In the analysis, each track is assigned a weight of $1/\varepsilon_{\mathrm{trk}}(\eta, p_{\mathrm{T}}, N_{\mathrm{pixel}})$ to account for track reconstruction effects.

4. Measurement method

The main experimental observable of this analysis is the mean transverse momentum $\langle p_{\mathrm{T}}\rangle$ of charged particles in an event as a function of N_{ch}, where $\langle p_{\mathrm{T}}\rangle$ and N_{ch} are measured within the same η and p_{T} ranges (otherwise, rapidity-dependent entropy fluctuations would lead to a reduced signal [17]). Charged particle p_{T} spectra for $p_{\mathrm{T}} > 0.3\,\mathrm{GeV}$ are measured for events in 50 GeV intervals of $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ from 3400 GeV to 5200 GeV, with tracking efficiency and misreconstruction effects corrected. To avoid any bias in estimating $\langle p_{\mathrm{T}}\rangle$ and N_{ch}, it is necessary to extrapolate the measured p_{T} spectra to the full p_{T} range. The resulting $\langle p_{\mathrm{T}}\rangle$ values (mean of the p_{T} spectra) from all $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ intervals are then plotted against the corresponding N_{ch} values (integral of the p_{T} spectra) to form the final observable. The $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ variable essentially serves as a centrality estimator to vary the initial medium entropy density and temperature. In particular, as the $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ values are obtained in a forward η range that does not overlap with the range used to measure the corresponding $\langle p_{\mathrm{T}}\rangle$ and N_{ch} values, potential biases are avoided. For example, hard processes originating early in the collision tend to fragment into large numbers of high-p_{T} particles, yet these particles may not reflect an increase in the entropy and temperature of the QGP medium.

The extrapolation of the p_{T} spectra to the full p_{T} range is performed by fitting a Hagedorn function [36] to the measured p_{T} spectra over the range of $0.4 < p_{\mathrm{T}} < 4.5\,\mathrm{GeV}$ in each $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ interval. This method is found to provide an excellent description of the data [37] and models (TRAJECTUM and HYDJET). The chosen p_{T} range for the fitting is varied to the evaluate corresponding uncertainties. The fitted functions are then used to extrapolate the missing portions of the p_{T} spectra in the low-p_{T} region.

As the extraction of the speed of sound mainly depends on the relative variation of $\langle p_{\mathrm{T}}\rangle$ with respect to N_{ch} (see equation (1)), normalized quantities, $\langle p_{\mathrm{T}}\rangle^{\mathrm{norm}} = \langle p_{\mathrm{T}}\rangle/\langle p_{\mathrm{T}}\rangle^{0}$ and $N_{\mathrm{ch}}^{\mathrm{norm}} = N_{\mathrm{ch}}/N_{\mathrm{ch}}^{0}$, are used as the primary observables, where the $\langle p_{\mathrm{T}}\rangle^{0}$ and N_{ch}^{0} represent the mean transverse

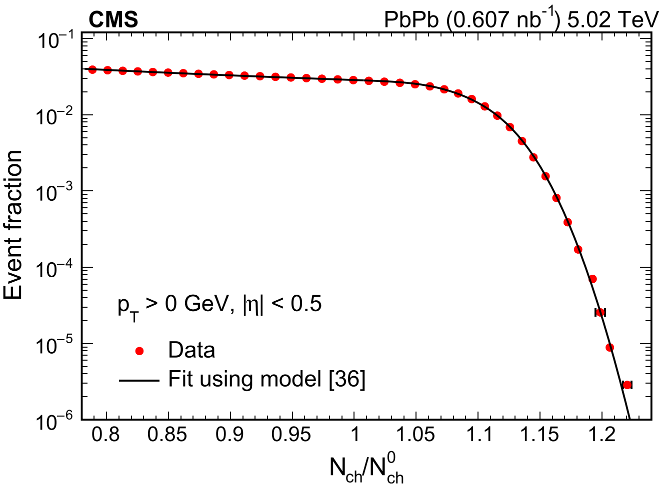

Figure 2. The event fraction distribution as a function of the charged-particle multiplicity, N_{ch}, within the kinematic range of $|\eta| < 0.5$ and extrapolated to the full p_{T} range, in PbPb collisions at $\sqrt{s_{\mathrm{NN}}} = 5.02\,\mathrm{TeV}$. The N_{ch} value is normalized by its value in the 0%–5% centrality class (N_{ch}^0). The curve represents a fit to the data using the Das *et al* model [39].

momentum and charged-particle multiplicity in a reference event class. Here, the centrality range chosen for the reference event class only needs to be close to that used for the speed of sound determination, and 5% most central events (as determined by $E_{\mathrm{T,\,sum}}^{\mathrm{HF}}$ and denoted '0%–5%') is used. By normalizing both $\langle p_{\mathrm{T}}\rangle$ and N_{ch} by their values in the reference event class, most of the systematic uncertainties can be minimized. The $\langle p_{\mathrm{T}}\rangle^0$ and N_{ch}^0 values obtained are found to be in good agreement with the ALICE results in the 0%–5% centrality range [37, 38]. Figure 2 shows the event fraction distribution as a function of the normalized multiplicity.

To extract the speed of sound, the expression that describes $\langle p_{\mathrm{T}}\rangle^{\mathrm{norm}}$ as a function of $N_{\mathrm{ch}}^{\mathrm{norm}}$ is taken from [17], as

$$\langle p_{\mathrm{T}}\rangle^{\mathrm{norm}} = \left(\frac{N_{\mathrm{ch}}^{\mathrm{norm}}}{\langle \overline{N_{\mathrm{ch}}^{\mathrm{knee}}} | N_{\mathrm{ch}}^{\mathrm{norm}}\rangle}\right)^{c_s^2}, \tag{2}$$

where,

$$\langle \overline{N_{\mathrm{ch}}^{\mathrm{knee}}} | N_{\mathrm{ch}}^{\mathrm{norm}}\rangle = N_{\mathrm{ch}}^{\mathrm{norm}} - \sigma\sqrt{\frac{2}{\pi}}\,\frac{\exp\left(-\frac{\left(N_{\mathrm{ch}}^{\mathrm{norm}} - \overline{N_{\mathrm{ch}}^{\mathrm{knee}}}\right)^2}{2\sigma^2}\right)}{\mathrm{erfc}\left(\frac{N_{\mathrm{ch}}^{\mathrm{norm}} - \overline{N_{\mathrm{ch}}^{\mathrm{knee}}}}{\sqrt{2}\sigma}\right)}. \tag{3}$$

Here, $\overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$ and σ represent the mean and root-mean-square width of the charged-particle multiplicity distribution at $b = 0$, normalized by N_{ch}^0. In figure 2, the $\overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$ value corresponds to the vicinity of the location beyond which the knee-shaped distribution starts rapidly falling. For the region of $N_{\mathrm{ch}}^{\mathrm{norm}} < \overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$, the $\langle \overline{N_{\mathrm{ch}}^{\mathrm{knee}}} | N_{\mathrm{ch}}^{\mathrm{norm}}\rangle$ variable approximately reduces to $N_{\mathrm{ch}}^{\mathrm{norm}}$, so equation (2) yields a value of unity. For the region of $N_{\mathrm{ch}}^{\mathrm{norm}} > \overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$, the $\langle \overline{N_{\mathrm{ch}}^{\mathrm{knee}}} | N_{\mathrm{ch}}^{\mathrm{norm}}\rangle$ variable saturates at $\overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$ for sufficiently large $N_{\mathrm{ch}}^{\mathrm{norm}}$. In this limit, equation (2) becomes a simple power function, with c_s^2 being the power of the function. The parameters $\overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$ and σ can be constrained by fitting the measured multiplicity distribution using the procedure described in [39]. The multiplicity distribution at fixed values of b is modeled using a Gaussian function. Integrating over b gives a minimum bias multiplicity distribution which can be fitted to data. As shown in figure 2, this fit provides a good description of the data. The results of this fit can be used to estimate the Gaussian mean and width at $b = 0$, yielding $\overline{N_{\mathrm{ch}}^{\mathrm{knee}}} = 1.11$ and $\sigma = 0.0272$ with negligible uncertainties. Using the extracted $\overline{N_{\mathrm{ch}}^{\mathrm{knee}}}$ and σ values, a fit to the measured $\langle p_{\mathrm{T}}\rangle^{\mathrm{norm}}$ as a function of $N_{\mathrm{ch}}^{\mathrm{norm}}$ is performed using equation (2), thereby extracting the speed of sound. In practice, we limit the fit to the very high-multiplicity region of $N_{\mathrm{ch}}^{\mathrm{norm}} > 1.14$, as will be discussed in detail later.

The dominant sources of systematic uncertainties for the measured $\langle p_{\mathrm{T}}\rangle^{\mathrm{norm}}$ and $N_{\mathrm{ch}}^{\mathrm{norm}}$ values originate from the tracking correction and the extrapolation to the full p_{T} range. As mentioned earlier, using normalized quantities minimizes the majority of the systematic uncertainties. Systematic uncertainties are directly evaluated for the normalized quantities, as well as for $\langle p_{\mathrm{T}}\rangle^0$ and c_s^2. The tracking correction uncertainty is evaluated by varying the default track selections to a set of looser or tighter values. The maximum deviation with respect to the default results is taken as a systematic uncertainty, which is found to be $\pm 0.01\,\mathrm{GeV}$ in $\langle p_{\mathrm{T}}\rangle^0$ and ± 0.002 in the fitted c_s^2 value. The p_{T} extrapolation uncertainty is estimated by varying the range of measured spectra fitted by the Hagedorn function to a lower limit of 0.3 or 0.5 GeV and an upper limit of 4 or 5 GeV. The resulting systematic uncertainty is found to be at most $\pm 0.023\,\mathrm{GeV}$ for $\langle p_{\mathrm{T}}\rangle^0$ and ± 0.012 for the c_s^2 value. Systematic uncertainties for c_s^2 associated with the choice of the lower fit limit in $N_{\mathrm{ch}}^{\mathrm{norm}}$ are estimated by varying the limit from 1.13 to 1.17, resulting in an uncertainty of ± 0.010 in c_s^2. Total uncertainties are obtained by adding the various sources in quadrature. Systematic uncertainties for $\langle p_{\mathrm{T}}\rangle^{\mathrm{norm}}$ are extracted point-by-point as a function of $N_{\mathrm{ch}}^{\mathrm{norm}}$.

5. Results

The observed multiplicity dependence of the average transverse momentum, both normalized by their values in the 0%–5% centrality class, is presented in figure 3, within the kinematic range of $|\eta| < 0.5$ and extrapolated to the full p_{T} range in central PbPb events. Hydrodynamic simulations from the TRAJECTUM [19, 40, 41] and Gardim *et al* [17] models are also shown for comparison. Both models use an EoS from lattice QCD calculations [42]. The TRAJECTUM model is a computational framework to simulate the full evolution of heavy ion collisions, which includes the modeling of initial stages, a viscous hydrodynamic phase with transport coefficients, and a hadronic gas phase. Parameters of the

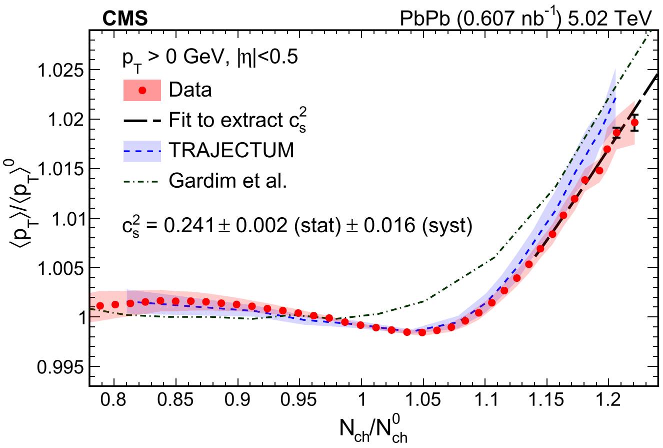

Figure 3. The average transverse momentum of charged particles, $\langle p_{\mathrm{T}} \rangle$, as a function of the charged-particle multiplicity, N_{ch}, within the kinematic range of $|\eta| < 0.5$ and extrapolated to the full p_{T} range in PbPb collisions at $\sqrt{s_{\mathrm{NN}}} = 5.02\,\mathrm{TeV}$. Both $\langle p_{\mathrm{T}} \rangle$ and N_{ch} are normalized by their values in the 0%–5% centrality class ($\langle p_{\mathrm{T}} \rangle^0$ and N_{ch}^0). Bars and the red band correspond to statistical and systematic uncertainties, respectively. Hydrodynamic simulations from the TRAJECTUM model [19] and the model by Gardim *et al* [17] are also shown for comparison. The dashed line is a fit to the data using equation (2) in the range of $N_{\mathrm{ch}}/N_{\mathrm{ch}}^0 > 1.14$.

TRAJECTUM model are constrained by a global Bayesian analysis of a variety of experimental observables [19], where the band shown corresponds to uncertainties within the allowed range of TRAJECTUM configuration parameters. The model of Gardim *et al* [17], besides the hydrodynamic phase, also considers the preequilibrium dynamics and hadronic interactions after thermal freeze-out. No uncertainties are evaluated for this model as only a single set of model parameters is used.

The $\langle p_{\mathrm{T}} \rangle^{\mathrm{norm}}$ value first shows a very weak declining trend toward a local minimum around $N_{\mathrm{ch}}^{\mathrm{norm}} \sim 1.05$. At higher multiplicities, corresponding to ultra-central PbPb events, a steep rise is observed, which is consistent with the expected increase in temperature with entropy density, as schematically illustrated in figure 1. The observed trend, including the minimum around $N_{\mathrm{ch}}^{\mathrm{norm}} \sim 1.05$, is qualitatively consistent with the prediction by the TRAJECTUM model. A slightly steeper rise at high multiplicities is observed for the TRAJECTUM simulation when compared with the data. This suggests that the speed of sound used in the model may be slightly larger than is found in the QGP. However, this difference is not significant within experimental and theoretical uncertainties. The model by Gardim *et al* also predicts a rise of $\langle p_{\mathrm{T}} \rangle^{\mathrm{norm}}$ at very high multiplicities, with a slope similar to that observed in the data. However, it shows a flat trend at lower multiplicities instead of the local minimum structure around $N_{\mathrm{ch}}^{\mathrm{norm}} \sim 1.05$ as seen in the data and the TRAJECTUM model. The origin of the observed local minimum is not currently understood.

To directly extract the speed of sound, the multiplicity dependence of the $\langle p_{\mathrm{T}} \rangle^{\mathrm{norm}}$ data in figure 3 is fitted by equation (2). Because the observed local minimum is not captured by the simplified model in equation (2), the fit is performed only in the high-multiplicity range with $N_{\mathrm{ch}}^{\mathrm{norm}} > 1.14$. The final result of the squared speed of sound is found to be $c_{\mathrm{s}}^2 = 0.241 \pm 0.002\,(\mathrm{stat}) \pm 0.016\,(\mathrm{syst})$ in natural units. The same fit is also performed to the prediction from the TRAJECTUM model, resulting in $c_{\mathrm{s}}^2 = 0.283 \pm 0.045$, where the model uncertainty is again determined within the allowed parameter space constrained by a global Bayesian analysis [19].

To constrain the EoS, a simultaneous> determination of c_{s}^2 and its corresponding temperature is necessary. Based on the hydrodynamic simulations discussed in [16, 17], the effective temperature (T_{eff}) of the QGP phase is found to be given approximately by $\langle p_{\mathrm{T}} \rangle/3$, with $T_{\mathrm{eff}} = \langle p_{\mathrm{T}} \rangle/3.07$ quoted [16] based on a soft EoS. While the scaling factor relating T_{eff} to $\langle p_{\mathrm{T}} \rangle$ can depend on specific model assumptions, the theoretical uncertainty in this value is believed to be small compared to the quoted experimental uncertainties, thereby having no impact on the main conclusions drawn in this paper. In essence, T_{eff} represents the initial temperature that a uniform fluid at rest would have if it possessed the same amount of energy and entropy as the QGP fluid does when it reaches its freeze-out state, the point at which the quarks become bound into hadrons. Due to longitudinal expansion and cooling, the T_{eff} value is generally lower than the initial temperature of the QGP fluid. Nevertheless, it still characterizes a temperature in the QGP phase, to which the extracted c_{s}^2 value based on the final-state $\langle p_{\mathrm{T}} \rangle$ and N_{ch} corresponds. Possible effects of shear and bulk viscosity are investigated in [16] and found to not impact this framework, as the shear viscosity increases $\langle p_{\mathrm{T}} \rangle$ by about the same amount that the bulk viscosity decreases it. The $\langle p_{\mathrm{T}} \rangle^0$ value is measured to be $658 \pm 25\,(\mathrm{syst})\,\mathrm{MeV}$, leading to a T_{eff} value for the ultra-central PbPb data of $219 \pm 8\,(\mathrm{syst})\,\mathrm{MeV}$ (it varies by at most 2% toward the very end of N_{ch} distribution within the 0%–5% centrality range). The statistical uncertainty is orders of magnitude smaller than the quoted systematic uncertainties.

Figure 4 depicts c_{s}^2 as a function of T_{eff}, with the CMS data point obtained from ultra-central PbPb collision data at $\sqrt{s_{\mathrm{NN}}} = 5.02\,\mathrm{TeV}$. The results are compared to the TRAJECTUM model, the c_{s}^2 value extracted in [16], and lattice QCD predictions of the c_{s}^2 value as a function of T [6]. The new CMS data allow for an unprecedented level of precision in the experimental determination of the speed of sound in an extended volume of QGP matter. The results exhibit excellent agreement with the lattice QCD prediction, with comparable uncertainties. Thus, our findings provide compelling and direct evidence for the formation of a deconfined QCD phase at LHC energies.

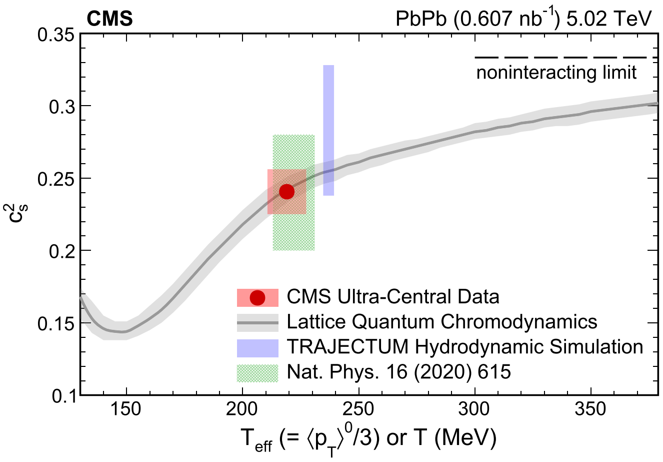

Figure 4. The speed of sound, c_s^2, as a function of the effective temperature, T_{eff}, with the CMS data point obtained from ultra-central PbPb collision data at $\sqrt{s_{\mathrm{NN}}} = 5.02\,\mathrm{TeV}$. The size of the red box indicates systematic uncertainties of c_s^2 and T_{eff}, while statistical uncertainties are smaller than the marker size. Values extracted from the Trajectum simulation [19] following the same fitting procedure as the data and from the earlier work [16] are presented as the other colored boxes. The curve shows the prediction of c_s^2 as a function of T from lattice QCD calculations [6]. The dashed line at the value of $1/3$ corresponds to the upper limit for noninteracting, massless gas ('ideal gas') systems [42].

6. Conclusion

In summary, this study presents a measurement with a new hydrodynamic probe in ultrarelativistic nuclear collisions that results in the most precise determination to date of the speed of sound in an extended volume of QGP matter. By determining the dependence of the average transverse momentum on the total multiplicity for charged particles in nearly head-on PbPb collisions at a center-of-mass energy per nucleon pair of 5.02 TeV, a squared speed of sound of $0.241 \pm 0.002\,(\mathrm{stat}) \pm 0.016\,(\mathrm{syst})$ in natural units is determined. The effective medium temperature, estimated using the mean transverse momentum, is $219 \pm 8\,(\mathrm{syst})\,\mathrm{MeV}$. The excellent agreement of lattice QCDs predictions with the experimental results provides strong evidence for the existence of a deconfined phase of matter at extremely high temperatures.

Data availability statement

Release and preservation of data used by the CMS Collaboration as the basis for publications is guided by the CMS policy as stated in CMS data preservation, re-use and open access policy.

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