

Search for the single production of doubly-charged Higgs bosons and constraints on their couplings from Bhabha scattering

OPAL Collaboration

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Abstract

A search for the single production of doubly-charged Higgs bosons is performed using e^+e^- collision data collected by the OPAL experiment at centre-of-mass energies between 189 GeV and 209 GeV. No evidence for the existence of $H^{\pm\pm}$ is observed. Upper limits are derived on h_{ee} , the Yukawa coupling of the $H^{\pm\pm}$ to like-signed electron pairs. A 95% confidence level upper limit of $h_{ee} < 0.071$ is inferred for $M(H^{\pm\pm}) < 160$ GeV assuming that the sum of the branching fractions of the $H^{\pm\pm}$ to all lepton flavour combinations is 100%. Additionally, indirect constraints on h_{ee} from Bhabha scattering at centre-of-mass energies between 183 GeV and 209 GeV, where the $H^{\pm\pm}$ would contribute via t -channel exchange, are derived for $M(H^{\pm\pm}) < 2$ TeV. These are the first results both from a single production search and on constraints from Bhabha scattering reported from LEP.

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1. Introduction

Some theories beyond the Standard Model predict the existence of doubly-charged Higgs bosons, $H^{\pm\pm}$,

including in left–right symmetric models [1], Higgs triplet models [2], and little Higgs models [3]. It has been particularly emphasized that a see-saw mechanism used to obtain light neutrinos in a model with heavy right-handed neutrinos can lead to a doubly-charged Higgs boson with a mass accessible to current and future colliders [4].

A review of experimental constraints on doubly-charged Higgs bosons is presented in [5]. The pair production of doubly-charged Higgs bosons has been considered in a previous OPAL publication [6], where masses less than 98.5 GeV are excluded for doubly-charged Higgs bosons in left–right symmetric models. DELPHI has obtained a limit of 97.3 GeV, independent of the lifetime of the $H^{\pm\pm}$ [7].

It has been noted that doubly-charged Higgs bosons may be singly produced in $e\gamma$ collisions, including in e^+e^- collisions where the γ is obtained from radiation from the other beam particle [8,9]. The diagrams for the direct production are shown in Fig. 1.

Doubly-charged Higgs bosons would decay into like-signed lepton or vector boson pairs, or to a W boson and a singly-charged Higgs boson. For masses less than twice the W boson mass, they would decay predominantly into like-signed leptons. Furthermore, in most models the WW branching fraction is negligible even for larger masses [9], therefore, the dominant decay mode, even for masses larger than twice the W boson mass, is the decay to like-signed leptons. Since the $H^{\pm\pm}$ naturally violates lepton number conservation, it can have mixed lepton flavour decay modes. Additionally, the Yukawa coupling of the $H^{\pm\pm}$ to the charged leptons $h_{\ell\ell}$ is model dependent, and is not generally determined directly by the lepton mass, so

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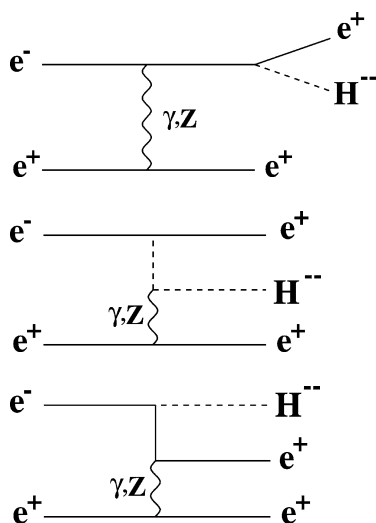


Fig. 1. Feynman diagrams contributing to the single production of H^{--} bosons in e^+e^- collisions. The three additional diagrams with “crossed” e^+ lines are not shown.

decays to all lepton flavour combinations need to be considered. It should be particularly noted that mixed lepton flavour decays are severely constrained by rare decay searches such as $\mu^+ \rightarrow e^+e^+e^-$ and $\mu \rightarrow e\gamma$.

In this Letter, we search for the single production of doubly-charged Higgs bosons, assuming the decays $H^{\pm\pm} \rightarrow \ell^\pm \ell'^\pm$ using 600.7 pb^{-1} of e^+e^- collision data with centre-of-mass energies $\sqrt{s} = 189\text{--}209 \text{ GeV}$ collected by the OPAL detector. Since the production cross-section depends only on h_{ee} , the Yukawa coupling of the $H^{\pm\pm}$ to like-signed electron pairs, the search is sensitive to this quantity.

We assume that the decay of a doubly-charged Higgs boson into a W boson and a singly-charged Higgs boson is negligible. We consider an $H^{\pm\pm}$ which couples to right-handed particles, but the results of the direct search quoted here are also valid for an $H^{\pm\pm}$ which couples only to left-handed particles [9]. All lepton flavour combinations are considered in the $H^{\pm\pm}$ decay ($ee, \mu\mu, \tau\tau, e\mu, e\tau, \mu\tau$). The lifetime of the $H^{\pm\pm}$ can be important, and in particular is non-negligible for $h_{\ell\ell} < 10^{-7}$; however, our search is not sensitive to such small Yukawa couplings.

A doubly-charged Higgs boson would also affect the Bhabha scattering cross-section via the t -channel exchange diagram shown in Fig. 2, causing a change in rate and in the observed angular distribution of the

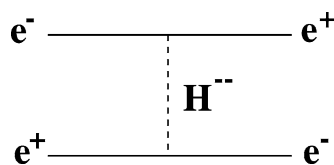


Fig. 2. Feynman diagram contributing to the process $e^+e^- \rightarrow e^+e^-$ due to doubly-charged Higgs boson t -channel exchange.

outgoing electron. Constraints have been derived for this process using data from lower energy colliders [5], but not previously from LEP.

In addition to the direct search results introduced above, we also derive indirect constraints on h_{ee} , the Yukawa coupling of $H^{\pm\pm}$ to electrons, using the differential cross-section of wide-angle Bhabha scattering measured by OPAL in 688.4 pb^{-1} of data collected at $\sqrt{s} = 183\text{--}209 \text{ GeV}$.

2. OPAL detector

The OPAL detector is described in detail in [10]. It is a multipurpose apparatus with almost complete solid angle coverage. The central detector consists of a silicon micro-strip detector and a system of gas-filled tracking chambers in a 0.435 T solenoidal magnetic field which is parallel to the beam axis. A lead-glass electromagnetic calorimeter with a presampler surrounds the central detector. In combination with the forward calorimeters, the forward scintillating-tile counters, and the silicon-tungsten luminometer, a geometrical acceptance is provided down to 25 mrad from the beam direction. The silicon-tungsten luminometer measures the integrated luminosity using small-angle Bhabha scattering events. The magnet return yoke is instrumented for hadron calorimetry, and is surrounded by several layers of muon chambers.

3. Direct search

3.1. Data samples and event simulation

The data samples are summarised in Table 1.

The process $e^+e^- \rightarrow e^\mp e^\mp H^{\pm\pm}$ is simulated with the PYTHIA6.150 [11] event generator. In the simulation, the Equivalent Photon Approximation (EPA) is used to give an effective flux of photons originating from the electrons or positrons. The upper limit

Table 1
Data samples used in the direct search analysis

E_{cm} (GeV)	$\langle E_{\text{cm}} \rangle$ (GeV)	$\int \mathcal{L}$ (pb $^{-1}$)
188–190	188.6	175.0
190–194	191.6	28.9
194–198	195.5	74.8
198–201	199.5	78.1
201–203	201.7	38.2
203–206	205.0	79.4
206–209	206.6	126.1
188–209	197.7	600.7

of the virtuality Q^2 of the photon is given by the scale of the hard scattering process.²⁴ The process $e^\pm\gamma \rightarrow e^\mp H^{\pm\pm}$ is simulated in a left–right symmetric model for an $H^{\pm\pm}$ which couples to right-handed particles using the calculations from [8]. The contribution from Z-exchange is negligible. In order to obtain the full signal cross-section, a cut which PYTHIA applies by default at a minimum of 1 GeV on the transverse momentum of the lepton which radiates the $H^{\pm\pm}$ is explicitly switched off. The cross-section and the angular distribution are checked with the calculations of [9], using COMPHEP [12].

Separate samples are simulated with the 6 different decay modes (ee , $\mu\mu$, $\tau\tau$, $e\mu$, $e\tau$, $\mu\tau$). Samples of 500 events each are generated for each of the average centre-of-mass energies listed in Table 1 for $H^{\pm\pm}$ masses in 5 GeV steps from 90–200 GeV. For masses larger than twice the W boson mass the decay $H^{\pm\pm} \rightarrow W^\pm W^\pm$ is kinematically allowed. Its partial width, however, is negligible in most models [9]. In this Letter, the branching fraction $\text{BR}(H^{\pm\pm} \rightarrow W^\pm W^\pm)$ is assumed to be zero.

The dominant Standard Model backgrounds in this analysis are from the four-fermion processes

$$e^+e^- \rightarrow \ell^+\ell^-\ell'^+\ell'^-,$$

including events from the so-called “multi-peripheral” diagrams,

$$e^+e^- \rightarrow e^+e^-\gamma^{(*)}\gamma^{(*)} \rightarrow e^+e^-\ell^+\ell^-,$$

and lepton pairs,

$$e^+e^- \rightarrow \ell^+\ell^-.$$

Four-fermion processes, except

$$e^+e^- \rightarrow e^+e^-\ell^+\ell^- \quad (\ell = e, \mu, \tau),$$

$$e^+e^- \rightarrow e^+e^-q\bar{q},$$

are simulated with the KORALW event generator [13]. The non-multi-peripheral part of the processes

$$e^+e^- \rightarrow e^+e^-\ell^+\ell^-,$$

$$e^+e^- \rightarrow e^+e^-q\bar{q},$$

is simulated with grc4f2.1 [14]. The multi-peripheral diagrams are simulated with the dedicated two-photon event generators Vermaseren [15] for

$$e^+e^- \rightarrow e^+e^-\gamma^{(*)}\gamma^{(*)} \rightarrow e^+e^-e^+e^-,$$

and BDK [16] for

$$e^+e^- \rightarrow e^+e^-\gamma^{(*)}\gamma^{(*)} \rightarrow e^+e^-\mu^+\mu^-,$$

$$e^+e^- \rightarrow e^+e^-\gamma^{(*)}\gamma^{(*)} \rightarrow e^+e^-\tau^+\tau^-.$$

The Monte Carlo generators PHOJET [17] (for $Q^2 < 4.5 \text{ GeV}^2$) and HERWIG [18] (for $Q^2 \geq 4.5 \text{ GeV}^2$) are used to simulate hadronic events from two-photon processes. Lepton pairs are simulated using the KK2f [19] generator for $\tau^+\tau^-(\gamma)$ and $\mu^+\mu^-(\gamma)$ events and NUNUGPV [20] for $\nu\bar{\nu}\gamma(\gamma)$. Bhabha scattering is simulated with BHWIDE [21] (when both the electron and positron scatter at least 12.5° from the beam axis) and TEEGG [22] (for the remaining phase space).

Multihadronic events, $q\bar{q}(\gamma)$, are simulated using KK2f [19]. RADCOR [23] is used to simulate multi-photon events from QED processes. They make a negligible contribution to the background.

Generated signal and background events are processed through the full simulation of the OPAL detector [24] and the same event analysis chain was applied to the simulated events as to the data.

3.2. Analysis

The signal final state consists of four charged leptons. Two like-sign leptons originate from the $H^{\pm\pm}$ decay and are expected to be visible in the detector in most cases. The electron or positron which originates from the $ee\gamma$ vertex (see Fig. 1) in general escapes through the beampipe. The electron or positron which originates from the $eeH^{\pm\pm}$ vertex is also forward

²⁴ Q^2 is the negative squared four-momentum transfer.

peaked; however, it enters the detector in a significant fraction of signal events. The analysis is therefore divided into a two-lepton and a three-lepton analysis. The final states in the three-lepton case contain three leptons visible in the detector, two of them have the same sign and could originate from the decay of a doubly-charged Higgs boson. In the two-lepton case, two like-signed leptons are required, as expected in the decay of a doubly-charged Higgs boson.

Leptons are identified as low multiplicity jets. Jets are reconstructed from charged particle tracks and energy deposits (clusters) in the electromagnetic and hadron calorimeters. Tracks and clusters are defined to be of “good” quality using the requirements of [25]. After the jet reconstruction, double-counting of energy between tracks and calorimeter clusters is corrected by reducing the calorimeter cluster energy by the expected energy deposition from associated charged tracks [25], including particle identification information.

No explicit electron or muon identification is required, since it is found that the jet-based analysis technique retains high efficiency while reducing the background to an acceptable level. The same analysis is used to search for all 6 possible lepton flavour combinations, and the results are valid for all leptonic decay modes of the $H^{\pm\pm}$. The final background is dominated by Standard Model processes containing four charged leptons. The analysis cuts are listed below. The cut values of the two-lepton and three-lepton analyses differ slightly.

The requirements for the two-lepton analysis are:

- (2.1) The preselection requires low multiplicity events [26]. The events are additionally required to have at least two and less than nine charged tracks. The sum of charged tracks and clusters in the electromagnetic calorimeter must be less than 16. Tracks and clusters are formed into jets using a cone algorithm [27] with a half-angle of 20 degrees and a minimum jet energy of 2.5 GeV, and it is required that there be exactly two jets with polar angles²⁵ satisfying $|\cos\theta| <$

0.95, and which are not precisely back-to-back (within 5°). Finally, the sum of the energies of the two jets reconstructed in the event must be greater than 20% of \sqrt{s} .

- (2.2) Ordering the jet energies by their magnitude ($E_{\text{jet1}} > E_{\text{jet2}}$), the following requirements are made:
 - (a) $E_{\text{jet1}} > 0.1\sqrt{s}$;
 - (b) $E_{\text{jet2}} > 0.05\sqrt{s}$;
 - (c) $E_{\text{jet1}} < 0.995E_{\text{beam}}$;
 - (d) $E_{\text{jet1}} + E_{\text{jet2}} < 0.95\sqrt{s}$.
- (2.3) The invariant mass M_{inv} of the two jets must satisfy $M_{\text{inv}} > 40$ GeV. Typical mass resolutions are about 4 GeV for ee and 10 GeV for $\mu\mu$. No mass reconstruction is possible for $\tau\tau$, due to the undetected neutrinos.
- (2.4) Bhabha scattering is rejected by requiring that the acollinearity angle, ϕ_{acol} , satisfies $\phi_{\text{acol}} > 25^\circ$. The angle ϕ_{acol} is defined to be 180° minus the opening angle of the two jets.
- (2.5) The polar angle of each jet must satisfy $|\cos\theta| < 0.75$. The $H^{\pm\pm}$ candidate jet polar angles are plotted in Fig. 3(a) and (b) after cuts (2.1)–(2.4).
- (2.6) Each jet associated to the $H^{\pm\pm}$ must have either one or three charged tracks. The number of charged tracks is plotted in Fig. 3(c) after cuts (2.1)–(2.5).
- (2.7) Defining the sum of the track charges within each jet as the “jet charge”, the product of the charges of the two jets must be equal to +1. This value is plotted in Fig. 3(d) after cuts (2.1)–(2.6).

The requirements for the three-lepton analysis are:

- (3.1) The preselection is identical to that in cut (2.1) except that exactly three reconstructed jets are required. The two jets which have the highest reconstructed mass, as described in cut (3.3), have to satisfy $|\cos\theta| < 0.95$ and must not be precisely back-to-back (within 5°). There is no $|\cos\theta|$ requirement for the third jet. Finally, the sum of the energies of the three jets reconstructed in the event must be greater than 20% of \sqrt{s} .

²⁵ OPAL uses a right-handed coordinate system where the $+z$ direction is along the electron beam and where $+x$ points to the centre of the LEP ring. The polar angle θ is defined with respect to the $+z$ direction and the azimuthal angle ϕ with respect to the $+x$

direction. The centre of the e^+e^- collision region defines the origin of the coordinate system.

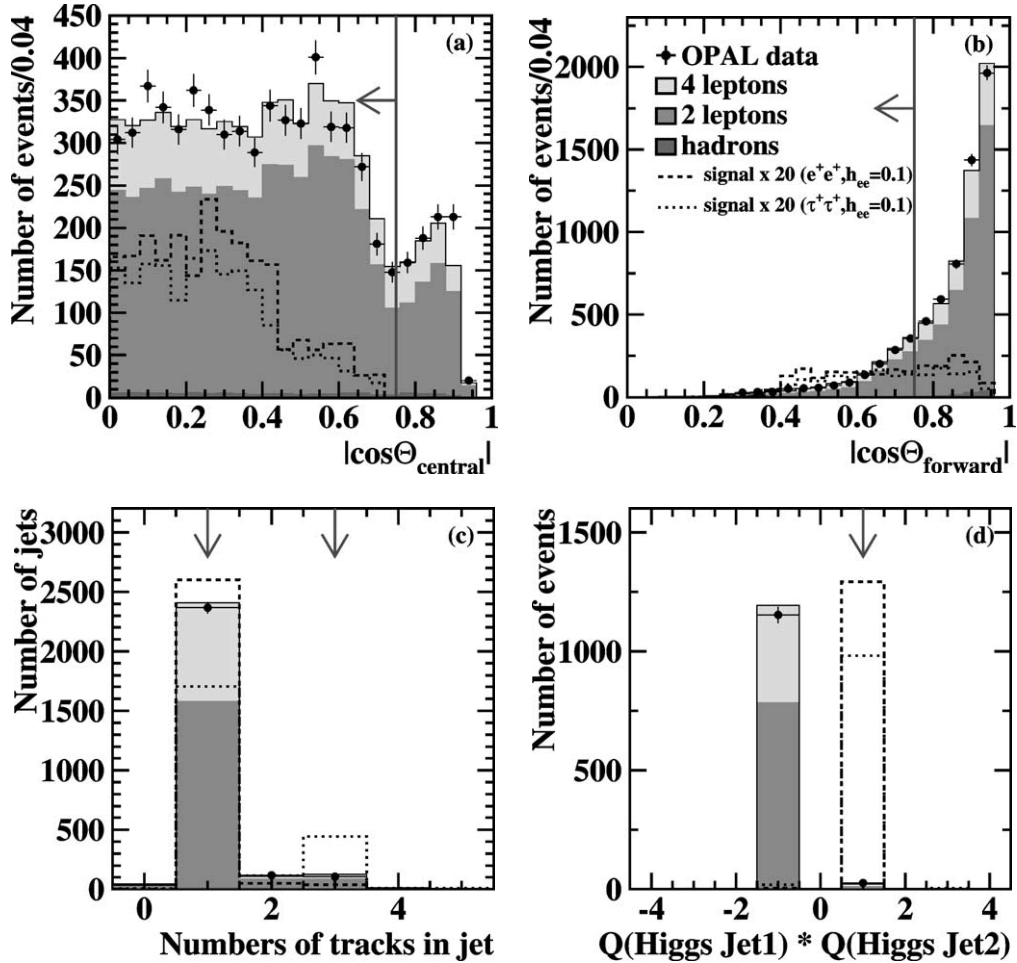


Fig. 3. Examples of some of the quantities used in the two-lepton analysis selection shown immediately before the corresponding cut is applied (see Section 3.2). The absolute values of the cosines of the polar angle of the more central and the more forward Higgs boson candidate jets are shown in (a) and (b), the number of charged tracks in each of the two $H^{\pm\pm}$ candidate jets in (c), and the product of the reconstructed charges of the two $H^{\pm\pm}$ candidate jets in (d). The points with error bars indicate the OPAL data and the shaded regions indicate the background expectation. Note that “hadrons” includes both $q\bar{q}(\gamma)$ and hadronic events from all 4-fermion processes. Two example signal expectations for a 130 GeV doubly-charged Higgs boson are also shown normalised to a cross-section corresponding to $h_{ee} = 0.1$ scaled by a factor 20, assuming either a 100% $H^{\pm\pm} \rightarrow ee$ branching ratio (dashed line) or a 100% $H^{\pm\pm} \rightarrow \tau\tau$ branching ratio (dotted line). The cut requirements are indicated by the arrows.

(3.2) Ordering the measured jet energies by their magnitude ($E_{\text{jet1}} > E_{\text{jet2}} > E_{\text{jet3}}$), the following requirements are made:

- (a) $E_{\text{jet1}} > 0.1\sqrt{s}$;
- (b) $E_{\text{jet2}} > 0.05\sqrt{s}$;
- (c) $E_{\text{jet3}} > 0.025\sqrt{s}$ or it must contain at least one good charged track;
- (d) $E_{\text{jet1}} < 0.995E_{\text{beam}}$;
- (e) $E_{\text{jet1}} + E_{\text{jet2}} + E_{\text{jet3}} < 0.95\sqrt{s}$.

(3.3) The jet energies are determined assuming that the measured jet direction is the same as the initial lepton direction for each of the reconstructed jets and that the missing electron or positron is recoiling along the beam axis. Using energy and momentum conservation to give four constraint equations, the four jet energies can be inferred (the lepton masses are neglected). Using this improved determination of the jet en-

ergies, the invariant masses are calculated for the three possible di-jet systems that can be constructed from the observed jets, and the two jets having the largest di-jet mass are considered as the $H^{\pm\pm}$ candidate jets with a “reconstructed Higgs boson mass” M_{rec} . The loss due to this assumption is negligible for $H^{\pm\pm}$ masses above 110 GeV, and is taken into account in the signal efficiency calculation. Since this search concentrates on the region above the mass limit from pair creation, it is further required that M_{rec} satisfy $M_{\text{rec}} > 80$ GeV. Typical mass resolutions are about 1 GeV for ee and $\mu\mu$ modes, and about 4 GeV for $\tau\tau$ decays. Note that in the latter case, no mass reconstruction from the jet energies would have been possible without this procedure, due to the undetected neutrinos.

- (3.4) Bhabha scattering is rejected by requiring that the acollinearity angle between the two $H^{\pm\pm}$ candidate jets satisfies $\phi_{\text{acol}} > 15^\circ$.
- (3.5) The polar angle of each jet associated to the $H^{\pm\pm}$ must satisfy $|\cos\theta| < 0.80$. The $H^{\pm\pm}$ candidate jet polar angles are plotted in Fig. 4(a) and (b) after cuts (3.1)–(3.4).
- (3.6) Each jet associated to the $H^{\pm\pm}$ must have either one or three charged tracks. The number of charged tracks is plotted in Fig. 4(c) after cuts (3.1)–(3.5).
- (3.7) Defining the sum of the track charges within each jet as the “jet charge”, the product of the charges of the two jets associated with the $H^{\pm\pm}$ must be equal to $+1$. This value is plotted in Fig. 4(d) after cuts (3.1)–(3.6).

The results are summarised in Table 2. The numbers of observed and expected events agree well after each cut in both analyses.

3.3. Systematic uncertainties

The largest background in the selection is from processes with four charged leptons in the final state, particularly from multi-peripheral “two-photon” processes. Of concern is the fact that, in our standard Monte Carlo background samples available at all centre-of-mass energies, the multi-peripheral diagrams are treated with specialised event generators which neglect interference with non-multi-peripheral

diagrams. Special samples of the full set of $e^+e^- \rightarrow e^+e^-\ell^+\ell^-$ diagrams, including interference, were prepared using grc4f2.2[14] at $\sqrt{s} = 206$ GeV to study this effect. The background using the full set of $e^+e^-\ell^+\ell^-$ diagrams including interference is in both analyses about 25% lower than our standard set of Monte Carlo generators. While grc4f2.2 includes interference effects, it has other differences with respect to our standard background simulations and cannot be used as the primary sample. We therefore simply assign a 25% systematic uncertainty on the $e^+e^-\ell^+\ell^-$ background according to this cross-check. Monte Carlo modelling of the variables used in the selection cuts can also induce systematic effects. The possible level of mismodelling is assessed by comparing data and background Monte Carlo for each variable after the preselection (cut (2.1) and (3.1), respectively) where the contribution from a signal would be negligible. Differences between the data and background Monte Carlo simulation are used to define a possible shift in each variable, and then the systematic uncertainties are evaluated by varying the cuts by these shifts. Both the final expected background and signal efficiencies are recalculated with these shifted cuts, and the full differences from the nominal values are assigned as systematic uncertainties.

The uncertainty of charge identification, used in cuts (2.7/3.7) in Section 3.2 to reject a significant fraction of the background, is estimated from a clean sample of Bhabha events selected by changing the cuts as follows. The cuts (2.2)(c) and (3.2)(d) are not applied. Cuts (2.2)(d) and (3.2)(e) are changed from $E_{\text{jet1}} + E_{\text{jet2}}(+E_{\text{jet3}}) < 0.95\sqrt{s}$ to $E_{\text{jet1}} + E_{\text{jet2}}(+E_{\text{jet3}}) > 0.95\sqrt{s}$. This sample consists mainly of Bhabha events and has no overlap with the search sample. The fraction of like-sign electron pairs is 2.0% in data and 1.7% in Monte Carlo. The systematic uncertainties on the background and signal efficiencies are evaluated by randomly changing the sign of the charge for 0.15% of the tracks, in order to increase the fraction of fake like-sign events by 0.3%, the observed difference between data and Monte Carlo in Bhabha sample. No significant $\cos\theta$ dependence of this difference was observed. An additional cross-check of τ pair events from calibration data taken at $\sqrt{s} = m_Z$ indicates that the charge-confusion systematic uncertainty is no larger than 0.15%. The full differences between the new background and efficiencies and

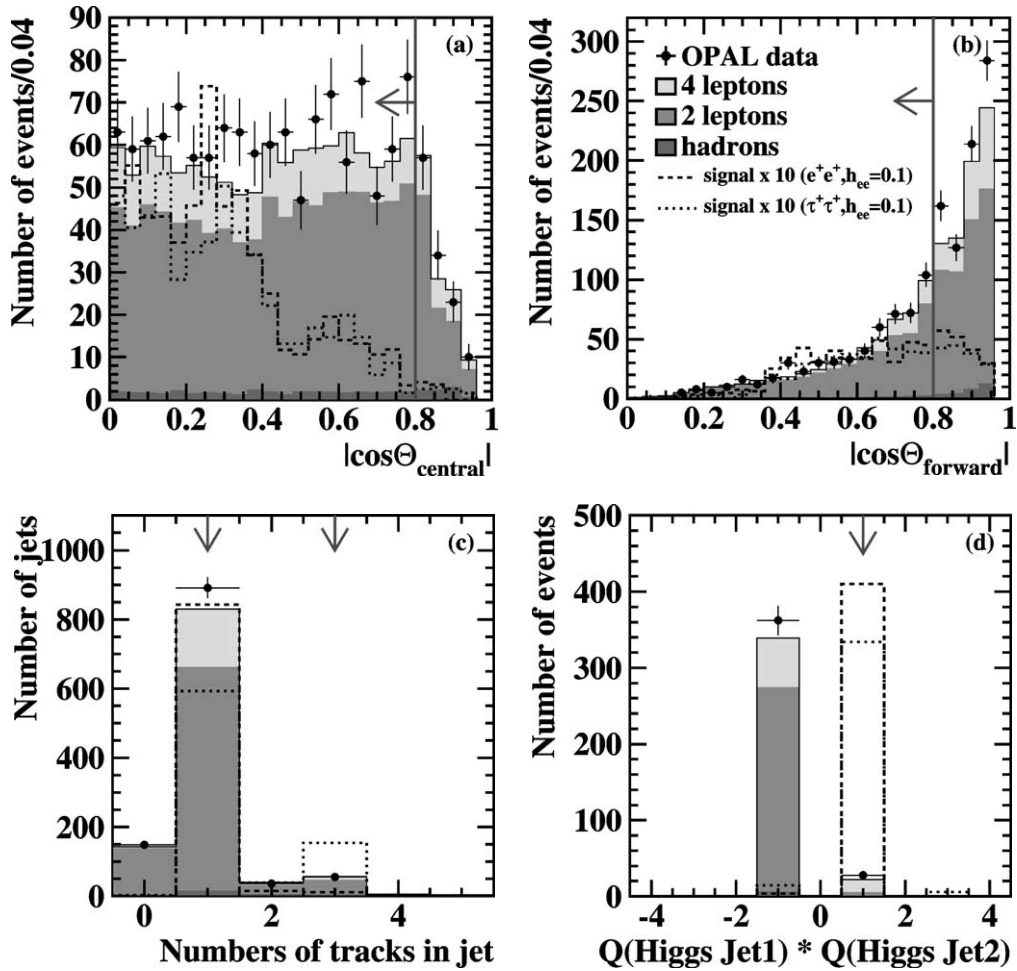


Fig. 4. Examples of some of the quantities used in the three-lepton analysis selection shown immediately before the corresponding cut is applied (see Section 3.2). The absolute values of the cosines of the polar angle of the more central and the more forward Higgs boson candidate jets are shown in (a) and (b), the number of charged tracks in each of the two $H^{\pm\pm}$ candidate jets in (c), and the product of the reconstructed charges of the two $H^{\pm\pm}$ candidate jets in (d). The points with error bars indicate the OPAL data and the shaded regions indicate the background expectation. Note that “hadrons” includes both $q\bar{q}(\gamma)$ and hadronic events from all 4-fermion processes. Two example signal expectations for a 130 GeV doubly-charged Higgs boson are also shown normalised to a cross-section corresponding to $h_{ee} = 0.1$ scaled by a factor 10, assuming either a 100% $H^{\pm\pm} \rightarrow ee$ branching ratio (dashed line) or a 100% $H^{\pm\pm} \rightarrow \tau\tau$ branching ratio (dotted line). The cut requirements are indicated by the arrows.

the nominal ones are taken as systematic uncertainties.

The systematic uncertainties are summarised in Table 3. Additional systematic uncertainties, such as on the integrated luminosity, are negligible.

3.4. Direct search results

In the two-lepton analysis the invariant mass of the two jets is calculated using the measured jet energies

and directions, because it is not possible to use the “angle-based” kinematic reconstruction described in Section 3.2 for the three-lepton analysis. The mass distribution is shown in Fig. 5 for events passing all cuts except the like-signed charge requirement (a), and also with all cuts applied (b). No excess of events which could imply the presence of a signal is observed in the data.

In the three-lepton analysis we calculate the $H^{\pm\pm}$ candidate reconstructed masses, M_{rec} , shown in Fig. 5,

Table 2

The number of remaining events in the data after each cut, and the number expected from Standard Model background sources. Also shown are the efficiencies of expected signal events for a 130 GeV doubly-charged Higgs boson assuming ee , $\mu\mu$ or $\tau\tau$ decays. The number of expected signal events for $h_{ee} = 0.1$ is shown in brackets assuming 100% branching ratio for the given decay mode. The errors due to Monte Carlo statistics are also listed for events surviving the full analysis

Cut	Data	Total Bkg.	$\ell^+\ell^-$	$4\text{-}\ell$	' $\gamma\gamma$ ' $e\ell\ell$	$q\bar{q}$	' $\gamma\gamma$ ' $e\bar{q}q$	Efficiency [%]		
								$e\bar{e}$	$\mu\mu$	$\tau\tau$
<i>Two-lepton analysis</i>										
(2.1)	19612	17659.3	13776.9	1249.6	2249.7	173.3	209.9	45.7	45.9	41.5
(2.2)	15168	14731.3	11381.3	1118.7	1971.5	158.1	101.8	44.3	39.1	36.4
(2.3)	13455	13002.6	10855.6	988.1	1026.5	120.8	11.6	44.3	39.0	35.0
(2.4)	6681	6685.9	5025.5	774.0	777.5	100.1	8.7	41.0	36.3	32.6
(2.5)	1318	1353.4	890.6	325.9	124.3	12.5	0.1	23.8	24.3	20.3
(2.6)	1181	1216.2	792.6	299.5	121.4	2.7	0.0	23.0	23.9	17.9
(2.7)	27	22.1	10.4	2.7	8.5	0.5	0.0	22.9	23.8	17.5
		± 1.7	± 1.3	± 0.2	± 1.1	± 0.1	± 0.0	± 1.9	± 2.0	± 2.0
								(64.6)	(67.3)	(49.1)
<i>Three-lepton analysis</i>										
(3.1)	40948	40899.7	7422.7	467.9	27011.1	260.1	5738.0	34.1	36.2	33.3
(3.2)	3203	2816.0	1685.9	153.3	778.9	63.1	134.8	22.7	24.0	21.1
(3.3)	2031	1912.0	1557.9	100.5	199.4	44.4	9.8	22.7	24.0	20.0
(3.4)	1359	1247.1	939.8	83.2	182.2	32.5	9.3	21.8	23.4	19.5
(3.5)	572	538.3	427.4	41.4	55.5	13.3	0.7	15.5	17.8	14.1
(3.6)	390	361.8	273.4	29.9	52.5	5.8	0.2	14.7	17.3	12.6
(3.7)	28	22.3	4.4	4.0	13.3	0.5	0.1	14.6	17.2	11.9
		± 1.6	± 0.7	± 0.3	± 1.4	± 0.1	± 0.0	± 2.0	± 2.0	± 2.1
								(41.0)	(48.8)	(33.4)
<i>Sum</i>										
\sum	55	44.4	14.8	6.8	21.8	1.0	0.1	37.5	41.0	29.3
		± 2.0	± 1.3	± 0.3	± 1.5	± 0.1	± 0.0	± 2.8	± 2.8	± 2.9
								(105.6)	(116.1)	(82.5)

Table 3

Systematic uncertainties on signal and background

Quantity	Variation	2-lepton analysis:		3-lepton analysis:	
		ΔBkg (%)	ΔSig (%)	ΔBkg (%)	ΔSig (%)
Jet $\cos\theta$	$\pm 0.5^\circ$	8	1	7	1
Jet energy	$\pm 1\%$	1	1	2	1
ϕ_{acol}	$\pm 0.5^\circ$	2	1	1	1
Charge misidentification	0.15%	14	1	4	1
Background modelling	(see text)	25	—	25	—
Monte Carlo statistics	—	8	10	7	14
Quadratic sum		31	10	27	14

using the “angle-based” kinematic reconstruction described in item (3.3) in Section 3.2. The mass distributions are shown both for events passing all cuts except

the like-signed charge requirement (c), and also with all cuts applied (d). Additionally, as a cross-check to ensure that no di-jet mass peak present after the event

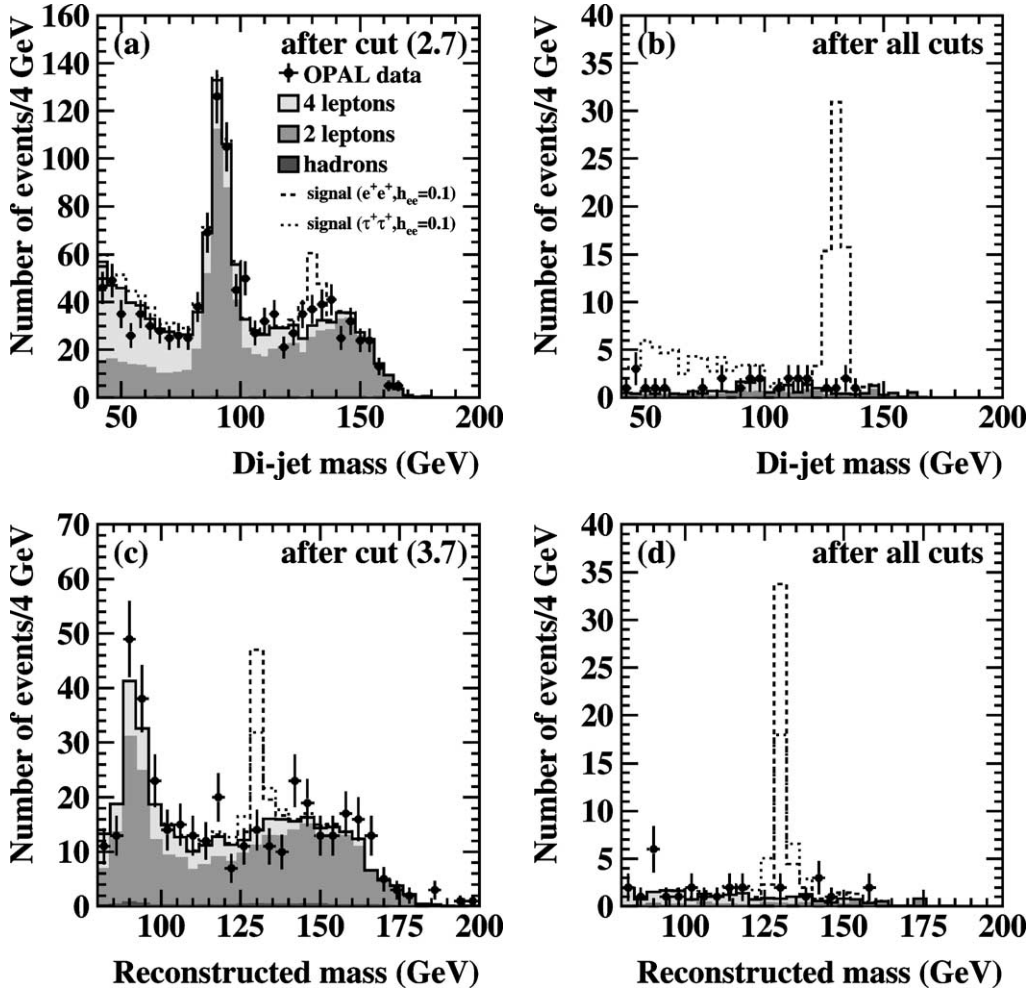


Fig. 5. The reconstructed $H^{\pm\pm}$ candidate mass distributions. The invariant di-jet mass is shown for the 2-lepton analysis both before and after the like-signed jet requirement (cut (2.7)) in (a) and (b), respectively. For the 3-lepton analysis, the reconstructed $H^{\pm\pm}$ mass using the jet angles as discussed in the text, is shown before and after cut (3.7) in (c) and (d), respectively. The points with error bars indicate the OPAL data and the shaded regions indicate the background expectation. Note that “hadrons” includes both $q\bar{q}(\gamma)$ and hadronic events from all 4-fermion processes. Two example signal expectations for a 130 GeV doubly-charged Higgs boson are also shown normalised to a cross-section corresponding to $h_{ee} = 0.1$, assuming either a 100% $H^{\pm\pm} \rightarrow ee$ branching ratio (dashed line) or a 100% $H^{\pm\pm} \rightarrow \tau\tau$ branching ratio (dotted line). Note that due to the undetected neutrinos from the tau-lepton decay there is no peak in the $H^{\pm\pm} \rightarrow \tau\tau$ signal sample of the 2-lepton analysis ((a) and (b)).

reconstruction is reduced by the angle-based method, the largest di-jet mass calculated from only the track and cluster information (Section 3.2) was examined. No excess of events which could imply the presence of a signal is observed in the data.

Limits are set on the $H^{\pm\pm}$ Yukawa coupling h_{ee} , assuming that the sum of the branching fractions of the $H^{\pm\pm}$ to all lepton flavour combinations is 100%. The efficiency for an arbitrary Higgs boson

mass is determined by linear interpolation between the simulated signal Monte Carlo samples. The number of observed events, together with the number of expected signal and background events from both the two-lepton and three-lepton analyses are combined using the likelihood ratio method described in [28], which incorporates the systematic uncertainties into the limits using a numerical convolution technique. For the purpose of extracting the limits, a ± 10 GeV

“sliding mass window” around the hypothetical Higgs boson mass is used. Events within this window are counted in data and Monte Carlo simulation. The hypothetical Higgs boson mass is varied in 1 GeV steps. The width of the mass window is chosen such that it contains most of the expected signal events. A small efficiency correction, typically around 5% for ee and $\mu\mu$ and 10% for $\tau\tau$, due to this window is applied. In the two-lepton analysis for any channel containing τ leptons no mass window cut is applied, because in this channel it is not possible to reconstruct the correct mass of the doubly-charged Higgs boson due to the undetected neutrinos.

The limits on h_{ee} are calculated using the efficiencies determined from the PYTHIA Monte Carlo samples and the production cross-sections are determined in a consistent manner using PYTHIA (see discussion in Section 3.1). No systematic uncertainty is assigned for theoretical uncertainties. The 95% confidence level limits on h_{ee} from combining both analyses are shown in Fig. 6(a)–(c) assuming a branching fraction of the doubly-charged Higgs boson into ee , $\mu\mu$, $\tau\tau$ of 100%, respectively. Strictly, due to the production mechanism involving non-zero h_{ee} , exactly 100% $\mu\mu$ or $\tau\tau$ decays are not possible, therefore, the latter limits should be considered for the case $h_{\mu\mu,\tau\tau} \gg h_{ee}$. In Fig. 6(d), for each mass the highest limit from all possible lepton flavour combinations is shown. An upper limit on $h_{ee} < 0.071$ is inferred for $M(H^{\pm\pm}) < 160$ GeV at the 95% confidence level, which is valid for all possible lepton flavour combinations in the decays. The limit is determined by the pure $\tau\tau$ case except for masses in excess of 170 GeV. For the case of pure ee decays the limit is $h_{ee} < 0.042$, and for $\mu\mu$ decays $h_{ee} < 0.049$, both for $M(H^{\pm\pm}) < 160$ GeV. For the mixed flavour decay modes $e\mu$, $e\tau$, and $\mu\tau$ the limit is between those for pure decays of the two involved flavours.

4. Indirect search

Doubly-charged Higgs bosons would contribute to Bhabha scattering via t -channel exchange as shown in Fig. 2. The Born level differential cross-section for Bhabha scattering including the exchange of a doubly-charged Higgs boson with right-handed couplings has been calculated in [5]. At high masses, $M(H^{\pm\pm}) \gg \sqrt{s}$, the cross-section is identical to that derived for

four-fermion contact interactions with right-handed currents [29] ($\eta_{RR} = 1$, $\eta_{LL} = \eta_{LR} = 0$), with the replacement of g/Λ by $h_{ee}/M(H^{\pm\pm})$ where h_{ee} is the Higgs coupling to electrons.²⁶ At values of $M(H^{\pm\pm})$ comparable to the centre-of-mass energy, this correspondence is modified by the inclusion of a propagator term. For comparison with the experimental data, QED radiative corrections are applied to the Born level terms for doubly-charged Higgs boson exchange and interference with Standard Model processes given in [5] using the program MIBA [30]. Initial state radiation is calculated up to $\mathcal{O}(\alpha^2)$ in the leading log approximation with soft photon exponentiation, and the $\mathcal{O}(\alpha)$ leading log final state QED correction is applied. The BHWIDE [21] program is used to calculate the Standard Model contribution to the differential cross-section. The theoretical predictions are calculated using the same acceptance cuts as are applied to the data.

This analysis uses OPAL measurements of the differential cross-section for $e^+e^- \rightarrow e^+e^-$ at centre-of-mass energies of 183–209 GeV [31,32]. The data between 203 GeV and 209 GeV are grouped into two sets with mean energies of approximately 205 GeV and 207 GeV. The total integrated luminosity of the data amounts to 688.4 pb^{-1} . These measurements cover the range $|\cos\theta| < 0.9$, in 15 bins of $\cos\theta$ (as defined in [32]), and correspond to $\theta_{\text{acol}} < 10^\circ$ where θ_{acol} is the acollinearity angle between electron and positron. It is verified that the effect of doubly-charged Higgs boson exchange on the low-angle Bhabha scattering cross-section has a negligible effect on the luminosity determination even for values of h_{ee} a few times larger than excluded by this measurement.

The measured differential cross-sections are fitted with the theoretical prediction using a χ^2 fit. The fit is performed for fixed values of the doubly-charged Higgs boson mass between 80 GeV and 2000 GeV, allowing the square of the coupling, h_{ee}^2 , to vary. Although only $h_{ee}^2 > 0$ is physically meaningful, in order to allow for the case where the data fluctuate in the opposite direction to that expected for doubly-charged Higgs boson exchange, both positive and negative values of h_{ee}^2 are allowed in the fit. Experimental and theoretical systematic uncertainties and their correlations

²⁶ In [5] h_{ee} is denoted g_{ee} .

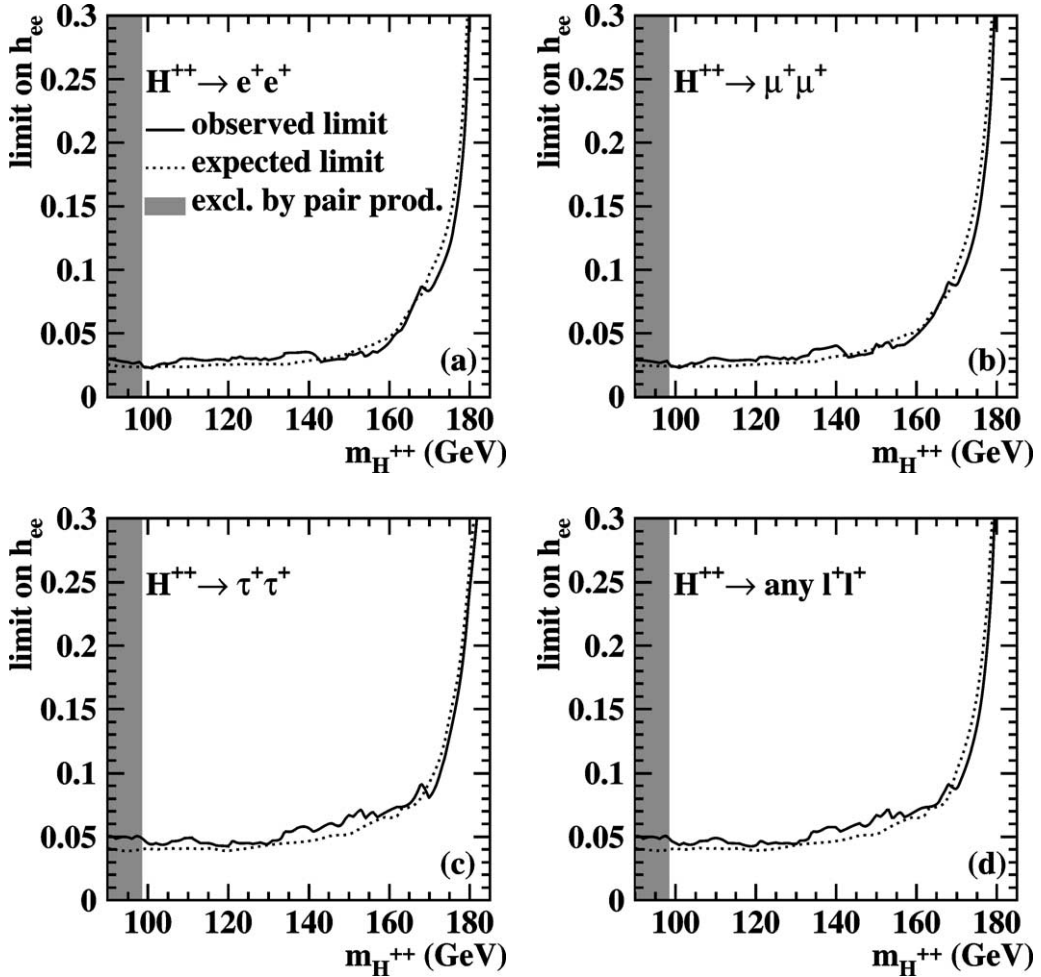


Fig. 6. Limits at the 95% confidence level on the Yukawa coupling h_{ee} assuming a 100% branching fraction of the $H^{\pm\pm}$ to (a) ee , (b) $\mu\mu$ and (c) $\tau\tau$. The limits are calculated with the combined results of the two-lepton and three-lepton analysis. In (b) and (c), the limits should be regarded as valid in the large branching fraction limit, since non-zero h_{ee} implies a non-zero electron branching fraction (see text). Since the ee and $\mu\mu$ efficiencies and mass resolutions are very similar, figures (a) and (b) are almost identical. The median expected limits assuming only Standard Model processes are shown by the dotted lines, while the actual limits inferred from the data are shown by the solid lines. In figure (d) the limit for arbitrary lepton flavour combinations (ee , $e\mu$, $e\tau$, $\mu\mu$, $\mu\tau$ and $\tau\tau$) is shown. It is determined by the pure $\tau\tau$ case except for masses in excess of 170 GeV. The shaded regions for masses below 98.5 GeV are excluded in left–right symmetric models by the OPAL pair production search [6].

are treated as discussed in [32]. The fitted values of h_{ee}^2 are consistent with zero for all masses, indicating that the data are consistent with the Standard Model prediction. For example, for a mass of 130 GeV the fitted value of h_{ee}^2 is 0.003 ± 0.011 , and the fit has a χ^2 of 97.0 for 119 degrees of freedom. Fig. 7 shows the ratio of the measured luminosity-weighted average differential cross-section at 183–207 GeV to the Standard

Model prediction, together with the results of the fit. 95% confidence level limits on the coupling as a function of mass were derived by integrating the likelihood function obtained from χ^2 over the region $h_{ee}^2 > 0$, and are shown in Fig. 8. The limits are considerably more stringent than those derived from PEP and PETRA data [5]. Fig. 9 shows the limits from the indirect search together with those from the direct search. The

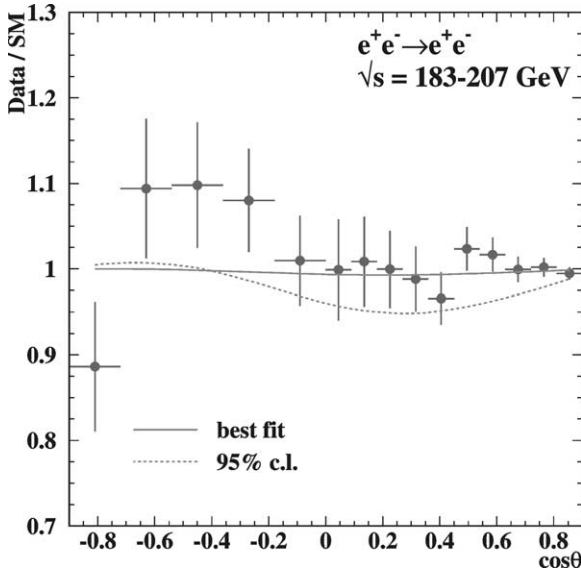


Fig. 7. Ratio of the measured luminosity-weighted average differential cross-section for $e^+e^- \rightarrow e^+e^-$ at 183–207 GeV to the Standard Model prediction. The points with error bars show the OPAL data, while the curves show theoretical predictions for a doubly-charged Higgs boson mass of 130 GeV. The solid curve corresponds to the best fit to all data, the dashed curve corresponds to a coupling equal to the 95% confidence level limit.

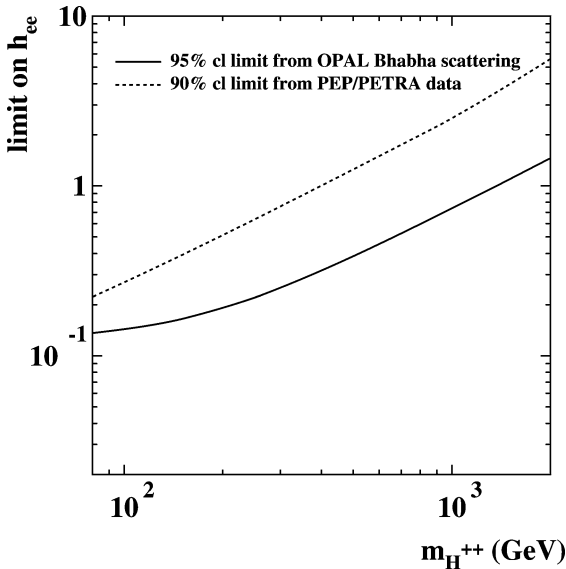


Fig. 8. Limits at the 95% confidence level on the Yukawa coupling h_{ee} as a function of $M(H^{\pm\pm})$ derived from Bhabha scattering data (solid line) for an $H^{\pm\pm}$ coupling to right-handed particles. Limits at 90% confidence level derived from PEP and PETRA data [5] are shown, as a dashed line, for comparison.

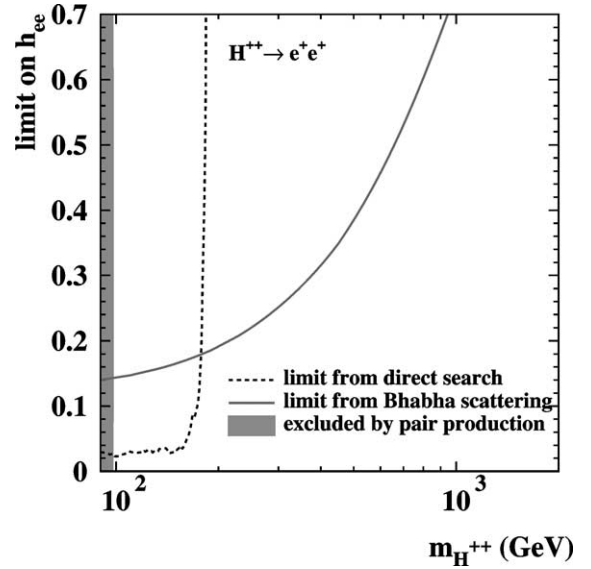


Fig. 9. Limits at the 95% confidence level on the Yukawa coupling h_{ee} assuming a 100% branching fraction of the $H^{\pm\pm} \rightarrow ee$. The direct limit is calculated with the combined results of the two-lepton and three-lepton analyses. The indirect limit on h_{ee} obtained from Bhabha scattering described in Section 4 is also shown. The shaded regions for masses below 98.5 GeV are excluded in left-right symmetric models by the OPAL pair production search [6].

indirect limits are less restrictive than those from the direct search at low masses, but extend to much higher masses.

5. Conclusion

A direct search for the single production of doubly-charged Higgs bosons has been performed. No evidence for the existence of $H^{\pm\pm}$ is observed. Upper limits are determined on the Higgs Yukawa coupling to like-signed electron pairs, h_{ee} . A 95% confidence level upper limit of $h_{ee} < 0.071$ is inferred for $M(H^{\pm\pm}) < 160$ GeV assuming that the sum of the branching fractions of the $H^{\pm\pm}$ to all lepton flavour combinations is 100%. Additionally, indirect constraints on h_{ee} for $M(H^{\pm\pm}) < 2$ TeV are derived from Bhabha scattering where the $H^{\pm\pm}$ would contribute via t -channel exchange for $M(H^{\pm\pm}) < 2$ TeV. These are the first results on both the single production search and constraints from Bhabha scattering reported from LEP.

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