

Excited states in the neutron-rich nucleus ^{25}F

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The structure of the nucleus $^{25}_{9}\text{F}$ was investigated through in-beam γ -ray spectroscopy of the fragmentation of ^{26}Ne and $^{27,28}\text{Na}$ ion beams. Based on the particle- γ and particle- $\gamma\gamma$ coincidence data, a level scheme was constructed and compared with shell model and coupled-cluster calculations. Some of the observed states were interpreted as quasi-single-particle states built on top of the closed-shell nucleus ^{24}O , while the others were described as states arising from coupling of a single proton to the 2^+ core excitation of ^{24}O .

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I. INTRODUCTION

Nuclei in the vicinity of doubly closed-shell nuclei are expected to exhibit simple shell model structures, which can often be interpreted in terms of selected single-particle degrees of freedom built on a closed-shell core. For stable nuclei, several such cases are known in the vicinity of, for example, nuclei like ^{16}O , ^{40}Ca , ^{56}Ni , ^{132}Sn , and ^{208}Pb that show such characteristics which are relatively well understood. However, more complex structures were observed for short-lived isotopes with extreme N/Z ratios.

Recent experiments indicated that $N = 16$ was a magic number close to the neutron drip line. This was first evidenced by neutron separation energy and cross section measurements [1]. Additional information came from the spectroscopy of the $N = 16$ isotones. The energy spectrum of ^{26}Ne resembled that of a vibrator nucleus [2], while only unbound excited states above the one-neutron spin density $S_n = 4.09(10)$ MeV [3], namely, a spin-parity $J^\pi = 2^+$ state at 4.7 MeV and a $J^\pi = 1^+$ at 5.3 MeV [4,5], were found in ^{24}O . Moreover, a small quadrupole deformation parameter, $\beta_2 = 0.15$ [5], as well as a large s wave spectroscopic factor of 1.74 ± 0.19 [6] was determined for the ground state of ^{24}O , providing further support for its closed $N = 16$ neutron shell. It was also proven that the drip line was reached at ^{24}O for the

oxygen isotopes reflected in the neutron separation energy step of 4.8 MeV between ^{24}O and ^{25}O [7]. A compatible value of 4.95 MeV was obtained from the $2J + 1$ (where J is the spin) weighted average energy of the two $J^\pi = 2^+$ and $J^\pi = 1^+$ states in ^{24}O [4].

Adding a single valence proton to the doubly closed-shell nucleus ^{24}O , ^{25}F is expected to have a rather simple structure. Its energy spectrum can be described up to the neutron separation energy of 4.36(12) MeV [8] as a few single-proton states coupled to the ground and first excited states of ^{24}O . Deviations from this straightforward picture may arise from the following: ^{25}F is expected to be located at the frontier of emerging new structures induced by intruder configurations, possibly leading to cluster configurations at low excitation energy. Recent cluster-model calculations [9] showed that the energy of cluster states associated with proton cross shell excitations increased with increasing neutron number among the fluorine isotopes. Conversely, the energy of a new class of intruder cluster configurations associated with coherent proton and neutron cross shell excitations decreased when approaching the neutron drip line. According to these calculations all the intruder states appeared above the neutron separation energy in ^{25}F [9]. The lowest-lying proton intruder state is predicted by shell model calculations using the Warburton-Brown psd-shell

(WBP) interaction in the *spsdpf* space at about 4.3 MeV in the ^{25}F nucleus [10,11]. Recently, two neutron-unbound states have been reported at 28(4) [10] and 300(170) [11] above the neutron separation energy which are expected to have intruder $J^\pi = 1/2^-$ and intruder $J^\pi = 3/2^-$ or normal $J^\pi = 5/2^+$ character, respectively.

The coupling of the weakly bound states to the continuum might also perturb the structure of ^{25}F . All the states arising from the coupling of a proton to the excited states of ^{24}O are expected to lie close to the neutron separation energy. The first excited state of ^{24}O itself is unbound by about 1 MeV [5]. The effects of continuum coupling were shown in the description of the properties of neutron-rich oxygen isotopes in Refs. [12,13]. Coupling to the nonresonant continuum and/or to eventual low-lying resonances may also play an important role in explaining the excitation energy spectrum of ^{25}F .

Preliminary papers on the study of ^{25}F from the fragmentation of the stable ^{36}S beam at the Grand Accelérateur National d'Ions Lourds (GANIL) were already published [14–16]. We reported on four γ lines, two of which were confirmed in an experiment performed at Rikagaku Kenkyusho (The Institute of Physical and Chemical Research) of Japan (RIKEN) [17]. In the final analysis of the spectra seven γ rays could be assigned to ^{25}F [18]; however, γ - γ coincidence was not available and only a tentative level scheme could be constructed for ^{25}F from these data. In the present paper we show the results obtained on ^{25}F at GANIL by a detailed in-beam γ -spectroscopic study via double-step fragmentation reaction. The experimental analyzes are interpreted in terms of both shell-model and coupled-cluster (CC) calculations. The effective Hilbert space for the shell-model calculations is defined by the $1s0d$ shell. CC calculations involve much larger effective Hilbert spaces, typically ten or more major oscillator shells. For states close to the separation energy, one expects that correlations from states in the continuum may play a larger role, suggesting, thereby, the need for larger Hilbert spaces. The larger dimensionalities mean, however, that only selected correlations are summed to infinite order in CC approaches, in contrast to configuration interaction calculations performed by the nuclear shell-model. For the latter, the many-nucleon eigenvalue problem is solved numerically exactly in a limited space. The degrees of freedom, of say the $1s0d$ shell studied here, may however not capture the relevant physics of more neutron-rich fluorine isotopes. These aspects will be discussed in our theoretical analysis of the experimental data.

II. EXPERIMENTAL METHODS

In the double-step fragmentation reaction a primary beam of ^{36}S at 77.5 MeV A with a mean intensity of 400 pnA was delivered by the two GANIL cyclotrons to induce fragmentation reactions in a carbon target of 348 mg/cm² thickness placed in the Superconducting Intense Source for Secondary Ions (SISSI) [19] device. The produced nuclei were selected by means of the ALPHA spectrometer equipped with a 130 mg/cm² Al wedge at the dispersive focal plane. The magnetic rigidity of the spectrometer and the optics of the beam line were optimized for the transmission of secondary beam particles with $N/Z \approx 5/3$ composed of ^{24}F , $^{25,26}\text{Ne}$,

$^{27,28}\text{Na}$, and $^{29,30}\text{Mg}$ nuclei with energies varying from 54 to 65 MeV A. These nuclei subsequently impinged on an “active” target, made of a plastic scintillator (103 mg/cm²) sandwiched by two carbon foils of 51 mg/cm² each, placed at the entrance of the SPEG spectrometer [20] to induce a secondary reaction. The plastic scintillator part of the active target was used to identify the incoming nuclei through energy loss and time-of-flight measurements. A total of 3×10^4 ^{25}F nuclei were produced in the fragmentation reaction, mainly from ^{26}Ne and $^{27,28}\text{Na}$ secondary beams. The ^{25}F nuclei were selected by the SPEG spectrometer and identified at its focal plane by a combined use of energy loss, total energy, time-of-flight, and focal-plane position information. This latter parameter served to correct the time-of-flight value from the various flight path lengths of the fragments in the spectrometer operated in a dispersive mode. The energy losses and positions of the fragments were determined by the use of an ionization (70 cm long) chamber and a set of two x - y drift chambers (2 cm thick and 80 cm wide each). The residual energy of the fragments was measured in a 2 cm thick plastic scintillator whose timing signal served to determine the time-of-flight T_{pl} with respect to the cyclotron radio frequency and to the plastic scintillator of the active target.

An array of 74 BaF₂ crystals was mounted in two hemispheres around the active target at a mean distance of 21 cm, covering a total solid angle of 80%. It was used to detect in-flight γ rays emitted with $v/c \simeq 0.34$ in coincidence with the ^{25}F fragments. The γ spectra were corrected for the Doppler-shift, yielding a full width at half maximum of about 12%. The detection limit was set below 100 keV for each detector in the laboratory frame. The efficiency of the BaF₂ array was determined by a GEANT4 simulation, which was calibrated using ^{137}Cs and ^{60}Co standard γ sources placed at the target position. A typical photopeak efficiency of about 20% is obtained at 1.3 MeV. Owing to the compact geometry of the array, high energy γ rays could scatter from one detector to another. To reduce the background and to achieve a higher efficiency at high energies, an add-back algorithm was used for the events where at least two neighboring detectors fired at the same time. For the proper determination of the efficiency, the source and simulated spectra were also analyzed using the add-back algorithm.

III. EXPERIMENTAL RESULTS

The Doppler-corrected γ -ray spectrum obtained for the ^{25}F nucleus is shown in Fig. 1. Two peaks can be seen at 750(30) and 1720(15) keV, while a broader structure including several γ lines is present between 3 and 4.5 MeV. This observation is in agreement with the results obtained in the fragmentation of the ^{36}S beam [14–16,18].

The decomposition of the broad structure into individual peaks requires the determination of the energy-dependent γ width of the peaks, which was obtained from single γ peaks observed in other reaction channels in the same experiment. An almost linear energy dependence of the peak width was observed. The response function of the BaF₂ crystals leads to a γ peak with a low-energy tail, the energy-dependent shape of which was simulated by use of the GEANT4 package. This

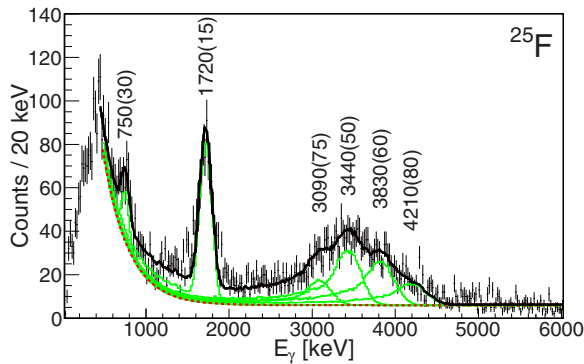


FIG. 1. (Color online) Decomposition of the γ -ray spectrum of ^{25}F . The solid black (dark) line shows the final fit which includes the response functions from GEANT4 simulation [green (light) solid curves] and the additional exponential background plotted as dashed red line.

low-energy tail is due to single and double escape as well as the Compton events which remain after the Compton suppression treatment. The line shape obtained with the simulation was tested successfully in the case of ^{22}O [3]. The fitting of the broad structure in the ^{25}F spectrum was made using the deduced γ line shape and an exponential background, yielding γ rays with energies of 3090(75), 3440(50), 3830(60), and 4210(80) keV as shown in Fig. 1. These energies are in agreement with the values deduced after reevaluation of the spectra obtained in the single step fragmentation [18].

For the level scheme construction, $\gamma\gamma$ -coincidence matrices were created. The most intense 1720 keV γ line was found to be in coincidence with itself, as well as with the 750 keV line. Furthermore, the 750 keV transition is in mutual coincidence with a 3440 keV γ ray as it can be seen in Fig. 2. In the analysis of the coincidence spectra the same line shapes and exponential background were assumed as for the single spectra.

The level scheme of ^{25}F was established by using $\gamma\gamma$ -coincidence matrices as well as energy and intensity balances. The energy and uncertainty of the excited states given below were determined taking into account the energies and uncertainties of all the γ rays connected to a given state via the fitting procedure of the RADWARE package [21].

The coincidence relation of the 1720 keV γ line with itself suggests two $\simeq 1720$ keV γ rays in cascade, establishing states at about 1720 and 3400 keV. As this energy overlaps with that of the 3440(50) keV transition within the experimental uncertainties, a level is proposed at 3440(21) keV. The γ line of 750 keV was chosen to feed the 3440 keV level since it was in coincidence with both γ rays of 1720 and 3440 keV. The sum $1720(15) + 1720(15) + 750(30) \simeq 4190(60)$ keV overlaps with the energy of the 4210(80) keV γ ray establishing a state at 4195(35) keV. The remaining transitions at 3090(75) and 3830(60) keV, which were not observed in coincidence with any transitions, were placed to directly feed the ground state, establishing levels at the corresponding energies. A weak 2140(30) keV line was observed in the single step fragmentation reaction [18], the intensity of which was enhanced in the multiplicity 2 spectrum, suggesting that it was in coincidence with another transition. Energetically such a

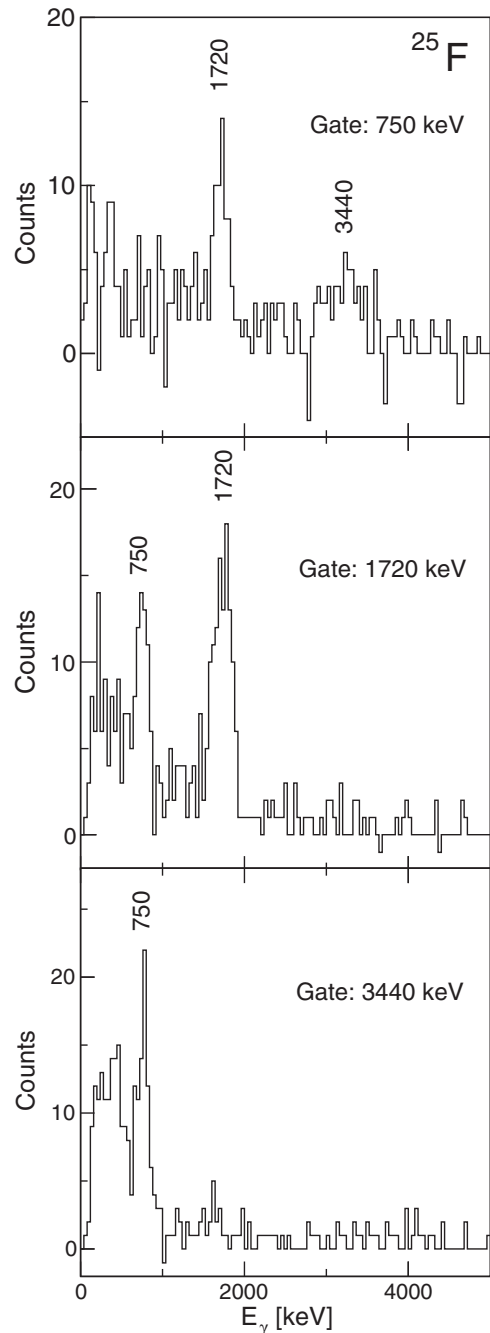


FIG. 2. Spectra of ^{25}F from $\gamma\gamma$ coincidence using the 750, 1720, and 3440 keV transitions as gates.

transition could connect the 1720 and 3830 keV states. In the single step reaction the 3830 keV line was the strongest, which explains why this transition was below the observation level in the present work.

The relative intensities of the 750(30), 3090(75), 3440(50), 3830(60), and 4210(80) keV transitions were determined on the basis of energy-dependent γ -ray efficiencies. The situation was more complex for the unresolved 1720 keV doublet, since the intensity of its members could not be derived directly. First the intensity ratio of the higher-lying member of the 1720 keV doublet and of the 3440 keV γ ray deexciting from

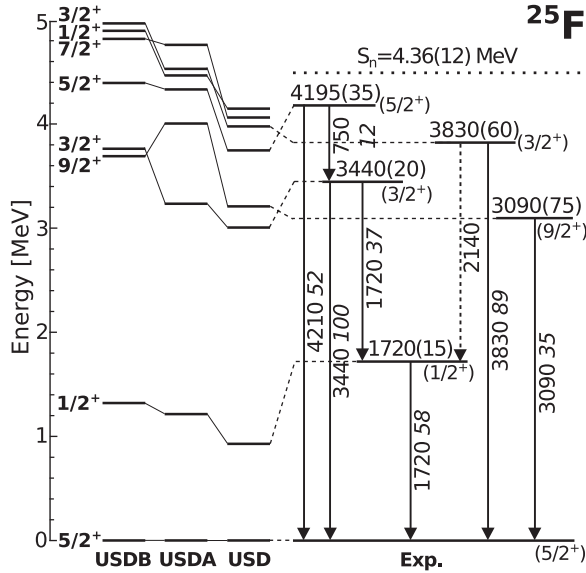


FIG. 3. Proposed level scheme of ^{25}F compared to the shell model calculations performed using the USD, USDA, and USDB interactions. Energies are given along the transitions as well as their relative intensities in italics. The uncertainties of the relative intensities are below 20%. The 2140 keV line has been adopted from Ref. [18].

the 3440(20) keV state was derived in a spectrum gated with the γ ray of 750 keV. From this ratio the intensity of the higher-lying member of the 1720 keV doublet was deduced using the corresponding peak areas in the single spectrum. The intensity of the lower-lying member of the doublet was obtained by subtracting the intensity of the higher lying component from the integral of the whole 1720 keV intensity. The resulting intensities of the γ transitions are presented in Fig. 3.

IV. DISCUSSION

The established level scheme of ^{25}F (Fig. 3) is compared to the results of shell-model calculations using effective interactions like the USD [22], USDA, and USDB [23] interactions defined for the $1s0d$ valence space. From a naive single-particle picture we expect that the ground state of ^{25}F has spin and parity $J^\pi = 5/2^+$. It corresponds to the filling of the $d_{5/2}$ single proton orbit above the ^{24}O core.

The first excited state at 1720(15) keV is supposed to correspond to the $J^\pi = 1/2^+$ state. In the simple single-particle picture, the $1/2^+$ state can be interpreted as a proton excitation to the $1s_{1/2}$ orbit. All other excited states, predicted above 3 MeV in the sd space, originate from the coupling of a proton single-particle state to excitations of the ^{24}O core in addition to the proton $d_{3/2}$ excitation. In addition, negative parity intruder or cluster states arising from a wider model space may also be present.

Considering that the state at 3440 keV decays both to the ground state of $J^\pi = (5/2^+)$ and the first excited $J^\pi = (1/2^+)$ state, a tentative $J^\pi = (3/2^+)$ assignment can be given to this state. Similarly, the 4195 keV state decays to the $J^\pi = (5/2^+)$ ground state and to the $J^\pi = (3/2^+)$ excited state, but not to the $J^\pi = (1/2^+)$ state, suggesting a $J^\pi = 5/2^+$ assignment

for this state. Although, the selection rules allow for spin $3/2$ or $7/2$, too, the missing transition to the $(1/2^+)$ state makes the spin $3/2$ assignment less probable, while the spin $7/2$ assignment is questioned by observation of the 750 keV transition, since a 750 keV $E2$ transition hardly could compete with an 4.2 MeV $M1$ one. With the above spin assignment the γ branching ratios are also in a qualitative agreement with the predictions of the sd shell model. While comparable intensities are expected for the transitions from the $J^\pi = 3/2_1^+$ state to the $J^\pi = 5/2_1^+$ and $J^\pi = 1/2_1^+$ ones, the $M1$ transition from the $J^\pi = 5/2_2^+$ to the $J^\pi = 5/2_1^+$ state is strongly hindered. This latter transition could only be observed since the E_γ^3 energy factor gives an enhancement of a factor of about 200.

All the shell-model interactions predict a $J^\pi = 9/2_1^+$ state below the $J^\pi = 5/2_2^+$ level. A good candidate for this $J^\pi = 9/2_1^+$ level is the level derived at 3090(75) keV. On the basis of shell model calculations and the fact that it directly decays to the ground state, the 3830 keV state's spin-parity can be limited to $1/2^+$ to $9/2^+$. Considering that the 2140(30) keV line observed in the single step reaction can also be assigned to the decay of the 3830 keV state, a spin $(3/2^+)$ might be assigned to this state. The branching ratios deduced from the single step reaction are consistent with the branching ratios calculated in the shell model for the $3/2_2^+$ state.

In Fig. 3, we can see that rather strong deviations can be found between the different versions of the shell-model interactions. The USD interaction fits reasonably well with the experimental data, while the USDA and USDB interactions stretch the energy spectrum too much. Assuming the above tentative spin assignments, the higher energy group of the experimental states lies about 500 keV below the predicted ones. This energy difference is similar to what was observed in ^{24}O . Since most of these states originate from the coupling of the proton single-particle states to the 2^+ and 1^+ states of ^{24}O the deviation might come from the inaccurate description of the ^{24}O core excitation. In this context we mention that these states of ^{24}O have a $\nu s_{1/2}d_{3/2}$ dominant configuration, where the $1s_{1/2}$ neutron is excited to the unbound $0d_{3/2}$ orbit. In this connection it may be interesting to mention that to describe the new data—among others—the neutron $d_{3/2}$ single-particle energy was increased by 330 and 440 keV in the USDA and USDB interactions, respectively, relative to the USD parametrization.

The unbound $d_{3/2}$ neutron may have an extended spatial distribution, which can result in a decrease of the interaction matrix elements in which the $d_{3/2}$ neutron is involved. Decreasing the $\nu s_{1/2}d_{3/2}$ interaction strength results in a decreased splitting of the $\nu s_{1/2}d_{3/2}$ doublet, yielding an increase in energy of the 2^+ state in ^{24}O . Decrease of the $\pi d_{5/2}\nu d_{3/2}$ interaction causes the decrease of the splitting of the $\pi d_{5/2} \oplus 2^+$ multiplet. As a result, both effects would worsen the agreement with experiment; thus the spatial extension of the unbound neutron $0d_{3/2}$ orbit cannot explain the observed deviation. A possible explanation may be the coupling to continuum configurations, as proposed in [12].

A low-energy bound intruder state with spin $J^\pi = 1/2^-$ was predicted around the neutron separation energy at 4296 keV using the $0p1s0d$ model space [10]. If bound, such a state would decay via an $E1$ transition, predominantly to the $J^\pi = 1/2^+$ state. Experimentally, we do not observe any state

with such decay property and energy, suggesting that the first intruder state in ^{25}F is unbound in agreement with the results of Franck *et al.* [10].

We also performed microscopic [24] CC calculations (see [25] for a textbook presentation) and compared them with experimental results. The coupled-cluster method is ideally suited to calculate properties of closed-shell nuclei and their immediate neighbors. While the shell-model restricts the model space to all possible determinants constructed from a small set of single-particle orbitals around a closed shell core, the CC method restricts the number of particle-hole excitations allowed in the determinant basis but uses a large single-particle space. Such particle-hole correlations are summed to infinite order depending on the level of approximation. Coupled-cluster theory is thus a nonperturbative method.

To describe ^{25}F , that has a proton attached to ^{24}O , we use particle-attached equation-of-motion coupled-cluster theory (PA-EOM-CC) [26,27]. In PA-EOM-CC the ^{25}F wave function is described by a linear expansion of 1p, 2p-1h, 3p-2h, ... excitations on top of the ground state of ^{24}O . In our earlier applications, we truncated the expansion at the 2p-1h level, which works particularly well for low-lying states that are dominated by 1p excitations from a closed-shell ground state [26–29].

In this work we are also interested in describing excited states in ^{25}F that can be viewed as a proton attached to the $J^\pi = 2^+$ excited state in ^{24}O . Clearly, to describe such states we need to go beyond the 2p-1h truncation level. As a rule of thumb, the level of approximation should be one order more than what is considered important for the state under study. We therefore include 3p-2h configurations as presented in Ref. [30], but as a first approach we only include the terms that determines the 3p-2h amplitudes in the coupled-cluster singles and doubles (CCSD) approximation. The similarity transformation generated by solving the CC equations induce additional many-body terms in the Hamiltonian. Only diagrams that contain at most two-body parts of the Hamiltonian have been included at the 3p-2h level, while all diagrams have been included up to the 2p-1h level.

We use interactions from chiral effective field theory at two different orders. First we look at the newly optimized chiral interaction at next-to-next-to-leading order (N^2LO) from Ref. [31]. Already at this order in the chiral expansion, three-body diagrams appear, but these have not been included here. This interaction resulted in an excellent agreement for both binding energies and selected excited states for oxygen isotopes. Second, we look at the chiral interaction at next-to-next-to-next-to-leading order (N^3LO) from Ref. [32]. Here, the effects of three-nucleon forces are treated as density-dependent corrections to the nucleon-nucleon interaction by integrating one nucleon in the leading-order chiral three-nucleon force over the Fermi sphere with a Fermi momentum k_F in symmetric nuclear matter [33]. We use the parameters already established in Ref. [29]. These parameters were used recently in a study of oxygen isotopes [29], and also for ^{26}F in Ref. [2].

Due to the increased computational cost of including 3p-2h configurations, the single-particle space is limited to a Hartree-Fock basis, built from $N = 10$ major harmonic oscillator shells. This is not large enough for the total binding

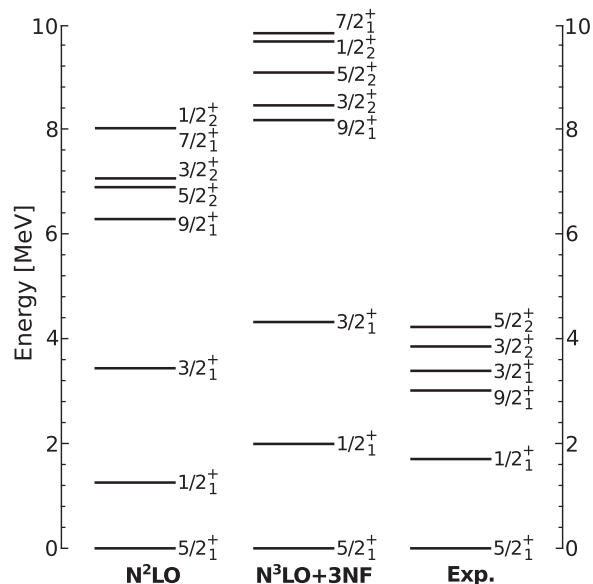


FIG. 4. Relative energies of ^{25}F calculated using coupled-cluster theory with two different interactions compared with experiment. The label N^2LO refers to the optimized chiral interaction at third order [31], while the interaction labeled $\text{N}^3\text{LO}+3\text{NF}$ is the chiral interaction of Ref. [32] where the effect of three nucleon forces are treated as a density-dependent nucleon-nucleon force (see text for details).

energies to be converged, but the relative energies exhibit a much faster convergence as function of the number of shells. The N^2LO interaction is rather soft, meaning that the relative energies are practically converged at this level. On the other hand our results using the N^3LO interaction exhibits a slower convergence rate, meaning that a larger model space is needed.

In Fig. 4 we plot the relative energies of selected positive parity states in ^{25}F calculated with the two different interactions, as well as experimental energies. The first excited $J^\pi = 1/2^+$ and $J^\pi = 3/2^+$ states are reproduced reasonably well by both interactions, but the remaining states are at too high excitation energies as compared to the data. This deviation might suggest that for these states, the current level of approximation is not good enough. Including 4p-3h configurations, where the latter account for 3p-3h excitations of the ^{24}O reference state, may further reduce the relative energies by 2 to 4 MeV. This guess is based on our observation that the relative energies for the high-lying states changes by 10 to 20 MeV by going from the 2p-1h to the 3p-2h truncation level. To get a better insight into which contributions are important, we examine here in more detail the structure of the calculated wave functions.

Table I lists the partial norms of the PA-EOM-CC wave functions calculated with the optimized N^2LO interaction. They sum up the parts of the wave function in 1p, 2p-1h, and 3p-2h configurations where the sum of all partial norms is one. We list only the partial norms for this N^2LO interaction since those obtained with the N^3LO interaction with density dependent three-body force produce qualitatively similar results.

TABLE I. Partial norms of wave functions using up to 3p-2h amplitudes on top of the ^{24}O reference state. The interaction is the newly optimized chiral interaction at third order (N^2LO) [31].

	1p	2p-1h	3p-2h
$5/2_1^+$	0.63	0.30	0.07
$1/2_1^+$	0.56	0.36	0.08
$9/2_1^+$	0.00	0.74	0.26
$3/2_1^+$	0.47	0.42	0.11
$3/2_2^+$	0.01	0.72	0.27
$5/2_2^+$	0.01	0.73	0.26
$1/2_2^+$	0.03	0.72	0.25
$7/2_1^+$	0.00	0.73	0.27

We see that the three presumed single-particle states have a large (30-40%) contribution from 1p-1h excitations of the ^{24}O reference state. If these were predominantly single-particle states, we would expect a contribution of approximately 10% or less from 2p-1h configurations. As it turns out, 2p-1h configurations are not enough to describe these states and they are not converged before 3p-2h configurations are included. We see that roughly 10% of the wave functions for these three states come from 3p-2h configurations, which implies a good level of convergence at this order. This is consistent with the shell-model calculations performed for these states. For the $J^\pi = 5/2^+$ ground state, the different USD interactions discussed above give wave functions with a 70 to 80% overlap with single-particle configurations, whereas the corresponding coupled-cluster numbers are 60 to 70%. For the first excited $J^\pi = 1/2^+$ and $J^\pi = 3/2^+$ states, the single-particle content of the wave functions has dropped significantly in both models.

For the three lowest lying positive parity states $J^\pi = 5/2_1^+$, $J^\pi = 3/2_1^+$, and $J^\pi = 1/2_1^+$, we note that the admixture from 3p-2h configurations does not exceed some 10%, meaning that we can interpret these states mainly in terms of 1p and 2p-1h configurations. The relative energies of these states are also well converged within the chosen set of configurations.

We notice also that for other positive parity states like the $J^\pi = 9/2_1^+$, $J^\pi = 3/2_2^+$, and $J^\pi = 5/2_2^+$ states, there are considerable admixtures from 3p-2h and more complicated configurations. Since additional 4p-3h configurations are expected to reduce the relative energies by 2 to 4 MeV, these changes are most likely larger than contributions that can be obtained from continuum effects. In Refs. [12,29] we found that the coupling to the continuum in ^{24}O gave a reduction in relative energy of 300 to 500 keV for states close to the neutron separation energy, just the energy difference found between the experimental results and the USDA/B calculations. As suggested by Table I, a suitable description of the high energy spectrum of ^{25}F would require the treatment of contributions above the (already large dimensionality) 3p-2h configurations. As this treatment is beyond the present computational limitations, a discussion of the continuum effects on these states has to be deferred to a later work.

We have also studied negative parity states within the PA-EOM-CC framework. We find a $J^\pi = 3/2^-$ state (not shown in Fig. 4) at approximately 10 MeV. It is identified as a spurious $J^\pi = 1^-$ center-of-mass excitation [34] built on

the $J^\pi = 5/2^+$ ground-state configuration and is therefore not considered a physical state [27,35]. At around 12 MeV we find a $J^\pi = 1/2^-$ state that is consistent with a physical state, but a large 3p-2h component suggests that it is not yet converged at this level of approximation.

V. CONCLUSION

We studied the structure of ^{25}F by the use of the in-beam γ spectroscopy technique from a double-step fragmentation reaction. Utilizing the γ -spectroscopic information, we constructed a level scheme for ^{25}F , including the states corresponding to the coupling to the core excitation of ^{24}O . Shell-model calculations using parametrized interactions for the $1s0d$ shell account for many of the general features of the energy spectrum, although a 500 keV deviation from the experiment was found in the USDA/B calculations for the temporarily assigned states of the $\pi d_{5/2} \oplus 2^+$ multiplet. All the states observed could be interpreted within the sd shell model; no bound negative parity intruder or cluster state has been found.

We also performed coupled-cluster calculations using recent interaction models from effective field theory. These calculations resulted in a very good agreement with experiment for the low-lying $J^\pi = 5/2^+$, $1/2^+$, and $3/2^+$ states. The analysis of the wave functions showed also a qualitatively similar picture in terms of single-particle states as the shell-model calculations. For higher-lying states that could be assigned to the $\pi d_{5/2} \oplus 2^+$ multiplet, our coupled-cluster calculations yielded excited states at too high excitation energies. An analysis of the structure of the wave functions indicated that correlations beyond the current truncation level might be necessary. In ^{24}O we found that the coupling to the continuum gave a reduction in relative energy of 300 to 500 keV for states close to the neutron separation energy, just the energy difference found between the experimental results and the USDA/B calculations. However, due to the large dimensionalities introduced by 3p-2h configurations in this work, we were not able to provide a proper estimate of continuum effects in ^{25}F .

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